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LAMINAR FORCED CONVECTION TO MIXTURES OF INERT GASES IN PARALLEL PLATE DUCTS

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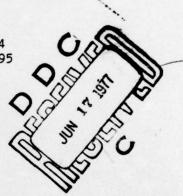
D. M. McEligot, M. F. Taylor and F. Durst

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D. M. McEligot 1,2, M. F. Taylor and F. Durst

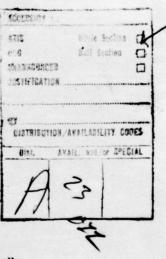
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ABSTRACT

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Mixtures of inert gases can be used to improve performance in closed gas turbine cycles. In the present work, heat transfer and wall friction parameters have been obtained numerically to demonstrate the effects of mixture composition and gas property variation for heating or cooling in regenerative heat exchangers of such cycles; the situation is modelled by laminar flow through short ducts with constant wall heat flux. For design predictions accounting for the effect of property variation, it is found that the property ratio method is better than the film temperature method for heat transfer while the latter method is preferable for apparent wall friction – with the proviso that the present definitions of the non-dimensional parameters be employed.

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NOMENCLATURE

	a,b,c	expanents for temperature-dependence of properties, equation (6)						
	c _p	specific heat at constant pressure						
	Dh	hydraulic diameter (twice plate spacing)						
	G .	average mass flux						
	g _c	dimensional constant						
	h	heat transfer coefficient						
	k	thermal conductivity						
	L	length						
	Ρ	pressure						
	P.9	exponents for property ratio method, equations (11) and (15)						
	T	absolute temperature						
	u .	axial velocity						
	v _b	bulk velocity						
	×	axial coordinate						
	Greek symbols							
	E/4,0	force constants in Lennard-Jones (6-12) potential						
	ц	absolute viscosity						
	P	density						
Dimensionless quantities								
	f, f _{ap}	apparent friction factor based on one-dimensional apparent wall shear stress, e.g., equation (13); f _a , lengthwise mean apparent friction factor						
	й	enthalpy, $(H - H_0)/(c_{pq}T_0)$						
	t*	length for velocity boundary layer, 4L/(DhRe)						

length for thermal boundary layer, $4L/(D_h RePr)$

Nusselt number, hDp/k

Nu

P	pressure, $2g_c \rho_o p/G^2$
Pr	Prandtl number, c µ/k
a ⁺	wall heat flux, q" DH/(koTo)
Re	Reynolds number, GD_{μ}/μ ; Re_{χ} , based on axial coordinate $PV_{bo}\chi/\mu$
, †	transverse velocity, (1/4)vRe/V _{bo}
× ⁺	axial coordinate for velocity boundary layer, 4 x/(DhRe)
×*	axial coordinate for thermal boundary layer, 4x/(DhRePr)
ÿ	transverse coordinate, y/D _h
Subscripts	
a	overall average or mean value (lengthwise)
ь	properties evaluated at bulk temperature
ср	based on constant property idealizations
•	in entry region,
f .	based on film temperature, $(T_w + J_b)/2$
fd	fully established or asymptotic value
m	lengthwise mean value
ref	reference
w .	properties evaluated at wall temperature
×	properties evaluated at local temperature
0	inlet conditions

The circumflex (^) represents non-dimensionalization with respect to the value of the quantity at the inlet, e.g., $\hat{\rho} = \rho/\rho_0$.

1. INTRODUCTION

The closed 3rayton cycle, or gas turbine cycle, has been suggested as an efficient, compact, versatile system for power plant and propulsion applications [1, 2]. With thermodynamic cycle studies, Bammert and Klein [3] showed that considerable savings in the total costs of a gas turbine cycle can be achieved by mixing helium with a heavier gas; the increase in density reduces the size of the turbomachines while the reduction in thermal conductivity increases the size of the heat exchangers so that an optimum occurs at an intermediate molecular weight. Approximate calculations indicate that when the heavier gas is another noble gas further improvements in heat transfer performance are possible, compared to pure gases at the same pressure, temperature and molecular weight [4]. However, the gaseous data and correlations for heat exchangers have been obtained with pure gases, primarily air, with negligible variation of the transport properties. Whether such results can be applied for mixtures of inert gases with temperature-dependent properties is a basic question which the present paper attacks.

In gas turbine cycles the regenerative heat exchanger is typically constructed of parallel plates, with short fins attached forming additional parallel surfaces. Consideration of the heat transfer performance versus pumping power requirements of these heat exchangers usually results in design for operation in the laminar or transitional flow regime. For laminar flow heat exchangers, the streamwise length of the fins is shortened in order to take advantage of the increased heat transfer coefficient of developing boundary layers by continually reinitiating the boundary layer. The thermal boundary condition is an approximately constant wall heat flux. As a guide to the effects of mixture composition and property variation in such geometries, the present paper – for a first objective – investigates the simultaneous development of laminar thermal and velocity boundary layers in the entry region of parallel plate ducts.

The following section briefly discusses related work which can be extended to provide improved guidance to the designer of regenerative heat exchangers for mixtures of noble gases. Section 3 summarizes pertinent knowledge of their transport properties and demonstrates the generalizations possible to reduce the analytical task. Since the consequent governing equations are nonlinear and coupled, they are solved numerically as outlined in section 4. The results of interest in design-lengthwise mean parameters – are presented in

section 5: first, the effect of composition at low heating rates and, then, the effects that the temperature-dependence of the transport properties cause on the heat transfer parameters and the wall friction parameters, separately. Finally, the major conclusions are reiterated in the last section.

2. PREVIOUS WORK

In preliminary design studies for closed gas turbine systems, Vanco [4] estimated relative heat transfer coefficients and pressure drop for binary mixtures of the inert gases with geometry held constant. For heat transfer he employed a constant property version of the Sieder-Tate relationship suggested by Kem [5]

This relation is essentially a thermal entry correlation of the Leveque form [6] which strictly applies only for a linear velocity profile as in the wall region of a fully established flow. With long tubes the Nusselt number should become constant rather than tending to zero as in the above relationship. For pressure drop calculations, Vanco used the friction factor for fully developed flow. His approach is still prevalent in industrial design of laminar flow heat exchangers.

As shown later, one of the main effects of varying mixture composition is to vary the Prandtl number. If the velocity profile is fully established before heating commences, the function $\operatorname{Nu_m}(L^{\pi})$ should depend an geometry and thermal boundary condition and be independent of Pr since the non-dimensional variables can be defined so that Pr does not appear in the governing energy equation or its boundary conditions [7]. Since thermal and shear boundary layers grow at different rates when the Prandtl number is not unity, the solutions to the simultaneous entry problem will vary with Prandtl number. In a numerical analysis Hwang and Fan [8] have shown this variation to be significant at low values of L^{π} for constant wall heat flux and constant fluid properties. For the comparable constant wall temperature problem, Schlünder [9] suggests applying the Pohlhausen solution in the immediate entry as

Num,
$$e = 0.664 (L^*)^{-1/2} Pr^{-1/6}$$
 (2)

in continuous functions of the form

$$Nu_{m} = \sqrt{Nu_{td}^{n} + Nu_{m,e}^{n}}$$
 (3)

for design computation.

Schade and McEligot [10] developed a numerical solution to examine the effect of air property variation for the simultaneous entry problem with strong heating or cooling applied to two parallel plates. They concluded that the local Nusselt number increases slightly and local friction factors increase severely with heating. To engineers accustomed to constant property analyses, the occurance of a large change in friction factor and pressure drop while the Nusselt number is changing only slightly might be surprising. Schade and McEligot showed that such comparisons are sensitive to the choice of the temperature at which the properties in the non-dimensional parameters are evaluated and that use of the constant property idealization can lead to either dangerous or conservative design, depending on the application.

For the heat exchanger applications, designers employ mean parameters and resort to empirical methods to correct for fluid property variation [11]. The two most common schemes are the reference temperature method and the property ratio method [12]. Based on experiment and approximate analysis, Kays recommends exponents for the latter method for various geometries and fluids. However, despite the ready availability of appropriate numerical methods for laminar flows for over a decade [13], these empirical correlations have not been refined and the question as to which is the more accurate method has not been answered. Accordingly, a second objective of the present paper is to examine this question and, if possible, to improve the exponents used in the property ratio method – for the parallel plate geometry and constant wall heat flux.

3. TRANSPORT PROPERTIES OF NOBLE GAS MIXTURES

While real gas properties can be employed in the numerical analysis in tabular or equation form directly, it is advantageous to exploit the similarities between different mixtures to reduce the number of computations necessary to cover the range of conditions of interest. Thus, if their behaviour can be generalized, fewer variations of parameters are required and the results are more concise and have greater usefulness. Accordingly, this section summarizes our knowledde of the pertinent transport properties. In conjunction with the following section it demonstrates that, as a first approximation, the variation in mixture composition can be reflected in the analysis in terms of a single variable parameter, the inlet Prandtl number.

The Lennard-Jones (6-12) potential can be employed in the Chapman-Enskog kinetic theory to predict thermal conductivity, viscosity and Prandtl number of binary mixtures of the inert gases [14]. There has been considerable experimental study of the pure gases, particularly helium and argon, but few data on their mixtures exist to check the predictions.

With force constants, E/K and O, as suggested by Hirschfelder, Curtiss and Bird [14] the prediction of the helium viscosity falls about eight percent below the data of Dawe and Smith [15] and Kalelkar and Keitin [16] at temperatures around 900°C. Likewise, the predicted thermal conductivity is about nine percent lower than the measurements of Saxena and Saxena [17] up to 1100°C. Similar discrepancies occur for xenon, but agreement with argon data is close. In general, agreement is good near room temperature for these gases.

Thornton [18] measured viscosity of the helium/xenon system at 20°C. His data agree with the predictions to within a few percent. On the other hand, the thermal conductivities obtained by Mason and von Ubish [19] at 520°C show an increasing divergence from the predicted curve as the fraction of helium is increased. Using force constants suggested by DiPippo and Kestin [20] for the helium component, instead of those of Hirschfelder, Curtiss and Bird, leads to essential agreement with the recommended value from the Thermophysical Properties Research Center [21].

Thermal conductivity and viscosity are presented against logarithmic coordinates in Figure 1 for helium, xenon, argon and some of their binary mixtures. These values are based on the Lennard-Jones (6-12) potential with the force constants $\mathcal{O}=2.158$, 3.292 and 4.055 A, $\mathcal{E}/\mathcal{K}=86.20$, 152.75 and 229 $^{\circ}$ K for helium, argon and xenon, respectively. The xenon values are from Hirschfelder, Curtiss and 3ird and the others from DiPippo and Kestin.

Viscosities of the noble gases and their mixtures differ only slightly with molecular weight (composition). The variation with temperature is approximately the same for all. Using an idealized temperature dependence,

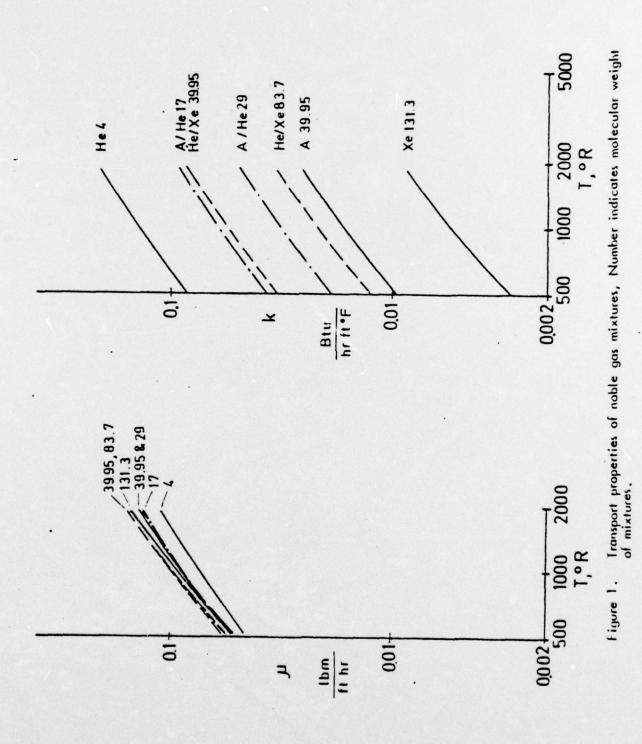
$$(\mu/\mu_{ref}) = (T/T_{ref})^{\alpha}$$
 (4a)

one finds the exponent "a" ranging from 0.7 to 0.8. On the other hand, the values of the thermal conductivity vary over an order of magnitude from xenon to helium. Again the temperature dependence is about the same for each and can be idealized as

$$(k/k_{ref}) = (T/T_{ref})^{b}$$
 (4b)

For the mixtures the exponent "b" is typically slightly less than "a" and ranges from about 0.7 to 0.75. As a first approximation, "a" and "b" could be taken as equal and the same value could be used for each of the mixtures. The specific heat is independent of temperature but varies with composition.

Perhaps the main surprise is the Prandtl number variation of the mixtures. Argon/helium and xenon/helium are shown in Figure 2; the other mixtures yield curves of the same shape. The temperature dependence is almost negligible since the power law exponents for k and μ differ so little. The Prandtl number of the pure gases is 2/3. However, each binary system shows a minimum at intermediate concentrations (molecular weight); for xenon/helium it is $Pr \approx 0.22$ at room temperature and is particularly broad. Other Prandtl number minima are: krypton/helium, 0.32; argon/helium, 0.42; krypton/neon, 0.53; xenon/argon, 0.57. Thus, Prandtl numbers in the range 0.5 to 2/3 can be obtained with several choices of binary mixture and concentration.



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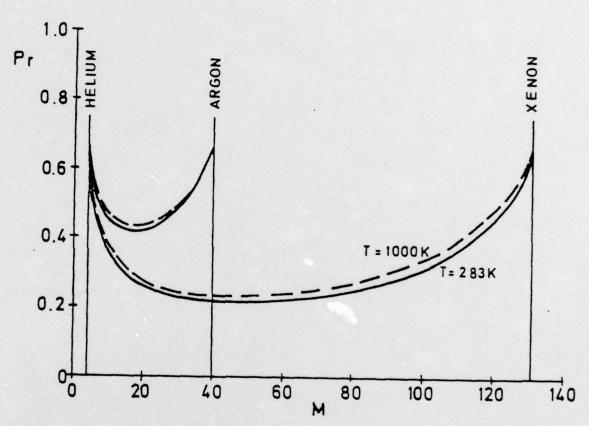


Figure 2. Prandtl number variation with concentration (molecular weight).

4. ANALYSIS

The details of the numerical analysis are straight forward and effectively the same as other methods for solution of coupled parabolic partial differential equations (e.g., reference [22]) so only the essentials will be outlined here. The number of parameters which must be varied to describe the range of conditions possible is determined by examination of the non-dimensional equations governing the problem.

Under the steady, internal boundary layer approximations – plus the assumptions that (a) mixture concentration remains constant, (b) Mach number << 1, (c) RePr > 100, and (d) natural convection is negligible – the governing equations can be written:

Continuity:
$$\frac{\partial \hat{\rho} \hat{u}}{\partial x} + \frac{\partial \hat{\rho} v^{+}}{\partial \overline{y}} = 0$$
 (5a)

$$\hat{\rho}\,\hat{u}\,\frac{\partial\hat{u}}{\partial x^*}+\hat{\rho}\,v^*\frac{\partial\hat{u}}{\partial \bar{y}}=-\frac{1}{2}\,\frac{d\bar{\rho}}{dx^*}+\frac{1}{4}\,\frac{\partial}{\partial\bar{y}}\,\left(\hat{\mu}\,\frac{\partial\hat{u}}{\partial\bar{y}}\right)$$

$$\hat{\rho} \hat{u} \frac{\partial \vec{H}}{\partial x^{+}} + \hat{\rho} v^{+} \frac{\partial \vec{H}}{\partial \bar{y}} = \frac{1}{4 \text{Pr}_{\bullet}} \frac{\partial}{\partial \bar{y}} \left(\hat{k} \frac{\partial \hat{T}}{\partial \bar{y}} \right)$$

Integral continuity:
$$\int_{0}^{\overline{y} \cdot 1/2} \hat{p} \, \hat{u} \, d \, \overline{y} = \frac{1}{2}$$
 (5d)

A circumflex (^) represents non-dimensionalization with respect to the value of the quantity at the entrance.

The gas properties may be idealized as:

$$\hat{p} = \frac{\hat{p}}{\hat{p}_{a} \hat{\uparrow}}; \quad \hat{\mu} = \hat{\uparrow}^{a}; \quad \hat{k} = \hat{\uparrow}^{b}; \quad \hat{e_{p}} = \hat{\uparrow}^{d}$$
 (6)

Initial conditions are: $\hat{T}(0,\bar{y}) = 1$, \bar{p}_0 specified and $\hat{u}(0,\bar{y}) = 1$, i.e.,

uniform entry. Boundary conditions are (1) the nonslip condition for velocities, (2) constant wall heat flux,

$$-\left(\hat{k}\cdot\frac{\partial\hat{T}}{\partial\bar{y}}\right)=Q^{+} \tag{7}$$

and (3) symmetry of the flow and boundary conditions with respect to the centre plane.

Examination of equations (5), (6) and (7) shows the free parameters of the mathematical statement are Pr_0 , Q^+ , \overline{p}_0 and the property exponents a,b and d. For the present paper a and b are taken as 0.75, d as zero and \overline{p}_0 is set sufficiently high that the Mach number is small in the range of interest. The parameters Pr_0 (corresponding to the mixture molecular weight) and Q^+ , the heating rate, remain variable.

The problem is solved numerically with program BAND, developed by Greif and McEligot [23] for flows between parallel plates with thermal radiative interaction using a band absorptance model. For the present calculations the capability to handle thermal radiation was suppressed, but property variation was included by choosing non-zero values of the exponents a and b. The numerical program is a finite control volume analysis using implicit algebraic equations to represent the governing equations; these equations are iterated at each axial step to treat their coupling and the nonlinear terms.

Mesh spacing increases in both the transverse and axial directions. For the results reported here, 81 transverse nodes were employed with the first usually at $(y/D_h) = 0.001$, and longitudinally there were ?0 steps per decade normally starting at $x_0 = 10^{-5}$. As noted by Worsoe-Schmidt and Leppert [13] the boundary layer approximations are not appropriate for $x^* < 10^{-3}$ so no results are reported for the initial two axial decades. Prior tests have shown convergence within 2 percent for Nu and within about 1 percent for f_{ap} with this grid.

5. RESULTS

Predictions have been obtained for the ranges $0.2 \le Pr \le 2/3$, corresponding to the Prandtl number variation from mixtures to pure gases, and $-2 \le Q^+ \le 100$. For $Q^+ \le 10$, \overline{p}_0 was taken as 10^3 giving an inlet Mach number of 0.035 while at higher Q^+ , $\overline{p}_0 = 10^4$ for $M_0 = 0.011$. Pertinent wall parameters are listed in Appendix A for all conditions studied in this investigation.

With heating at a constant wall heat flux, the local bulk temperature increases continuously; with constant specific heat this increase is linear. The viscosity increases with temperature, so the local bulk Reynolds number (GD/Mbx) decreases in the axial direction and the flow would be expected to remain laminar. Density is inversely proportional to temperature so the flow accelerates as it is heated, also normally a stabilizing influence. This continuous acceleration prohibits the occurance of invariant velocity and temperature profiles. The wall temperature is larger than the bulk temperature in order to transfer energy to the gas so, at a given cross section, the viscosity and thermal conductivity will be higher near the wall and the density will be lower. Consequently, parameters defined in terms of wall properties will have different values than those using bulk properties in their definition. While the increase in viscosity is expected to increase the wall shear stress and decrease the velocity near the wall, thus increasing thermal resistance, the increase in thermal conductivity is a factor tending to reduce thermal resistance. Likewise, the expansion due to reduced density counteracts the decrease due to viscosity to some extent. Near the entrance, the increased viscosity and decreased density at the wall augment the non-slip condition causing transverse flow away from the wall; this flow also carries thermal energy. Further downstream the transverse flow decreases. (These various effects are reversed with cooling). While the individual effects of these phenomena can be forecast in same cases, their combined effect is not obvious and may vary with geometry. Since density, viscosity and thermal conductivity variations are all of the same order of magnitude, none dominates in such a way that a single-parameter, closed-form analysis would appear feasible.

For the present study the presentation of results emphasizes mean parameters which are useful to design engineers; however, local Nusselt numbers and friction factors are included in the Tables for those interested. With pro-

perty variation, parameters can take a number of values depending on how the temperatures which are used to evaluate the pertinent properties are defined. To avoid later confusion in the use of these predictions it is necessary to be explicit in the definitions.

Local bulk temperature, $T_{\rm bx}$, is the temperature corresponding to the total enthalpy flow at a cross section, i.e., the socalled mixing cup temperature. Average bulk temperature, $T_{\rm ba}$, is the arithmetic average of the local value at length L and the inlet value

$$T_{ba} = (T_{ba} + T_{bl})/2$$

While an integral average is normally used for the average wall temperature in analyses it is not of use to the designer who lacks the detailed knowledge of $T_w(x)$; accordingly, we define T_{wa} as an arithmetic average

, which becomes, for constant wall heat flux,

$$T_{wa} = (T_{bo} + T_{wL})/2$$

(Thus, T_{wa} is lower than $\int T_{w}(x) dx/L$). Then, in agreement with normal practise, an average film temperature is chosen as

$$T_{fa} = (T_{wa} + T_{ba})/2$$

Comparable subscripting is used for indentifying the temperature at which properties are evaluated, e.g., $\mu_{ba} = \mu(T_{ba})$, and similarly for non-dimensional parameters, $Re_{ba} = GD_h/\mu_{ba}$. Average pressure is also taken at the arithmetic average of inlet and exit, L. Other quantities will be defined later, as used.

Predictions for constant fluid properties

For small heating rates or low temperature differences the constant properties idealization is often adequate for practical applications. For example, one sees in Figure 1 that a 50° C difference in temperature causes a change in thermal conductivity of about 17% at room temperature and 4% at 1000° K. In Figure 3 are plotted the mean Nusselt numbers versus non-dimensional length, $L^{\bullet}=4L/(D_{h}RePr)$, as obtained by setting the property variation exponents, a

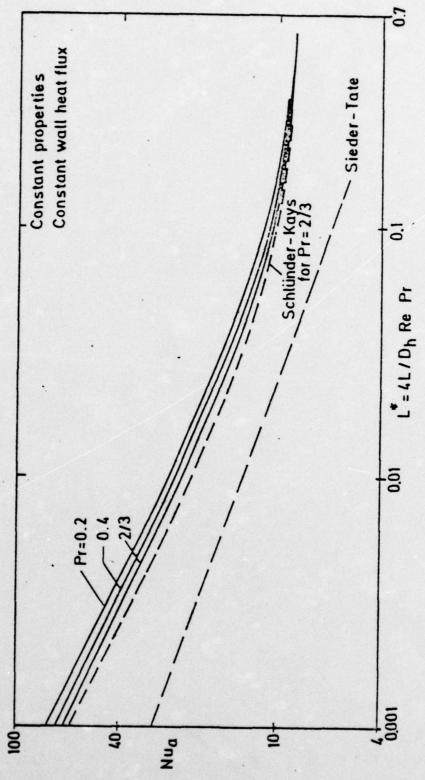


Figure 3. Moan heat transfer predictions under constant properties Idealization.

and b, at zero and holding $\hat{\rho}=1$. The mean Nusselt number is defined in terms of the integral average heat transfer coefficient as

$$Nu_a = \frac{ha \cdot Dh}{k} = \frac{Dh}{kL} \int_{0}^{L} h(x) dx$$

From Figure 3 one sees that in the entry $Nu_a(L^*)$ increases as the Prandtl number is reduced by employing different inert gas mixtures. This increase is approximately 17% as Pr varies from 2/3 to 0.2. Since, for constant properties, the development of the shear boundary layer is a function of $x^+ = 4x/(D_h Re)$ and is independent of Prandtl number (i.e., it is a solution of equations (5a,b and d), the velocity profile is more fully developed for Pr = 2/3 than for Pr = 0.2 at any specified value of L^* . Consequently, the increase in Nu_a coincides with a higher average velocity gradient, $(\partial u/\partial y)_a$, between 0 and L^* ; with increased velocities near the wall an increase in heat transfer parameters is reasonable.

Also plotted on Figure 3 is the Sieder-Tate correlation (1) which has been employed in practise for heat transfer calculations in comparable situations. In addition to differing by as much as a factor of two in the range of interest, this calculation fails to show the proper trend with L* except approximately in the limited range $0.01 < L^* < 0.1$. At L* ≈ 0.001 the mean Nusselt number varies as $(L^*)^{-0.47}$ rather than $(L^*)^{-1/3}$ as in the Sieder-Tate correlation. Failure of this correlation to account for Prandtl number dependence has been mentioned earlier.

Based on an integral/superposition method Kays [12] suggests that the local Nusselt number for the simultaneous growth of laminar external boundary layers can be predicted by

$$Nu_{x} = 0.453 \text{ Pr}^{1/3} \text{ Re}_{x}^{1/2}$$
 (8)

for the thermal boundary condition of a constant wall heat flux. For the immediate entry where $T_{bx} \approx T_{o}$ and $V_{bx} \approx V_{o}$, this relation can be transformed to a mean Nusselt number

$$Nu_{m,e} = 0.906 Pr^{1/3} Re_{D_h}^{1/2} D_h^{1/2} L^{1/2}$$
 (9)

In this case, equation (3), the form suggested by Schlünder [9], would become

$$Nu_{a} = \left[8.235^{2} + 1.812^{2}/(P_{r}^{1/3}L^{2})\right]^{1/2}$$
 (10)

This equation is shown on Figure 3 for a pure gas, Pr = 2/3, and can be seen to agree with the numerical prediction within about ten percent. For the mixtures agreement is as good or better, the comparisons are shown individually later.

It should be reemphasized that while the increase in $Nu(L^*)$ is of the order of 15 percent, when replacing a pure gas by a xenon-helium mixture, the gain in heat transfer coefficient is much greater. For example, taking equation (9) as an approximation for short fins in conjunction with Figures 1 and 2, one can see that for a given geometry and Reynolds number the effect of replacing pure argon by a helium/xenon mixture is to increase the heat transfer coefficient about 2.4 times.

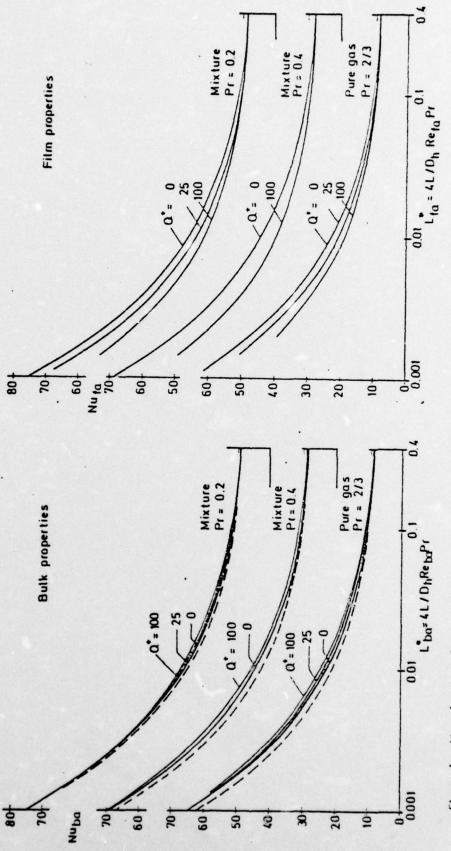
Predictions of heat transfer with property variation

When fluid properties vary significantly due to high heating rates and related large temperature variations in the flow field, the numerical values of non-dimensional parameters such as Nu and Re depend on the temperature at which their properties are evaluated. Two methods of accounting for the property variation are common in practise: the film temperature approach and the property-ratio approach [12]. In the film temperature approach it is assumed that using properties evaluated at T_f will allow direct use of predictions obtained under the constant property idealization. In the property ratio approach the bulk temperature is used for the properties and the effect of heating is represented as

$$Nu/Nu_{cp} = (T_w/T_b)^p \tag{11}$$

for gases. The effectiveness of these two methods for heat transfer predictions is examined in the present section; wall friction predictions are considered in the following section.

Figure 4 demonstrates the apparent effect of heating rate on the mean Nusselt number when the properties are taken at T_{ba} , i.e., the property ratio approach. Thus, L_{ba} is defined as $4L/(D_hRe_{ba}Fr)=4L k_{ba}/(c_pGD_h^2)$. The effect of heating the gas in a slight increase in Nu_{ba} compared to the prediction of Nu for constant properties $(Q^+ \rightarrow 0)$ at the same value of L_{ba} . For a



Princesson I

Figure 5. Mean heat transfer predictions in tems of film properties. Figure 4. Mean heat transfer predictions in terms of bulk properties. Dashed line represents equation (10).

given value of Q^+ , Nu_{ba} is increased more for the pure gases than for the mixtures, but this result is partially a consequence of the lower wall-to-bulk temperature reached by the gas with the lower Prandtl number. However, for $Q^+ = 100$, which leads to $(T_{wa}/T_{ba})_{max}$ around 2.5 to 2.7, the increase is only of the order of five percent. Thus, the exponent p in equation (11) would be about 0.05, which can probably be considered negligible for most practical cases. On the other hand, for $Q^+ = 100$ and Pr = 2/3 the local Nusselt number Nu_{bx} reaches a peak increase over the constant property prediction of about twenty-five percent at a location where $T_{wx}/T_{bx} \approx 3$ so the mean Nusselt number is affected even less than the local value.

To examine whether the film temperature approach is an improvement, the results are plotted in Figure 5 with properties evaluated at the average film temperature, $Nu_{fa} = h_a D_h/k_{fa}$ and $L_{fa}^* = 4Lk_{fa}/c_p GD_h^2$. With this choice the mean Nusselt numbers decrease as the heating rate is raised. This effect is of approximately the same magnitude for the pure gases as for the mixtures with their lower Prandtl numbers. The most significant abservation is that the reduction is about three times as great as the change when using bulk properties (Figure 4) so there is no advantage in accuracy with the film temperature approach.

When heating the gas T_w and T_b both increase, T_w more rapidly in the entrance region and then both at approximately the same rate at larger distances. Consequently, T_w/T_b first increases and then approaches unity axially. Likewise, $Nu_{ba}(L^*)$ is first greater than $Nu_{cp}(L^*)$ and then approaches $Nu_{cp}(L^*)$ downstream. In this situation correlations of the form of equation (11) make sense. In contrast, when cooling the gas – as on the apposite side of a regenerative heat exchanger – both temperatures decrease continuously such that, in the limit, T_w/T_b approaches zero or T_b/T_w approaches infinity. The relative property variation across the flow becomes greater downstream instead of less as in the heating situation. This limit is not likely to be of concern in a gas turbine cycle since the lowest temperature is set by the inlet temperature of the side to be heated, but might be a difficulty in cryogenic applications. In the cooling range covered in the scope of the present paper, $Nu_{ba}(L_{ba}^*)$ remained sufficiently close to $Nu_{cp}(L^*)$ to neglect the difference.

If one calls equation (10) the Schlünder-Kays relation and plots it as

a dashed line on each of the subfigures of Figure 4, it may be seen that agreement with the constant property result is good for each case but successively better as Pr is reduced. Since $Nu_{ba}(L_{ba}^*)$ is predicted so closely by $Nu_{cp}(L^*)$ it is worthwhile to improve the equation so it may be used by the designer over the range of interest. By adjusting the Prandtl number dependence of $Nu_{m,e}$ for better prediction at $L^* = 0.001$, one may obtain

$$Nu_{a} = \left[8.235^{2} + 1.931^{2}/(P_{r}^{0.254}L^{*})\right]^{1/2}$$
 (12)

This equation still is of the order of five percent lower than the numerical prediction near $L^{\#}=0.01$. For moderate heating rates it would be within about ten percent of the numerical results and would be low; in heat exchanger design this would lead to units slightly longer than necessary. Alternatively, the constants in equation (12) could be optimized for another range at the expense of the accuracy of predictions in the immediate entry.

Prediction of wall friction with property variation

While most analyses presently available for developing flows present friction results in terms of the wall shear stress evaluated from the velocity gradient at the wall (f_s) , this approach is not of use to the designer when the velocity profile is changing substantially, as in the entry or when a gas is heated. To predict the required pressure drop with a one-dimensional design procedure, one uses the "apparent" friction factor, f_{ap} , based on the wall shear determined by treating the momentum change as one-dimensional. The same treatment is often employed in experiments where size prohibits velocity profile measurements. Both methods of presentation can be chosen with numerical results; consequently, Bankston and McEligot [22] were able to demonstrate (a) the numerical values of f_s and f_a can differ substantially and (b) discrepancies earlier thought to exist between experiments and analyses were primarily due to the differences in the definition used for the friction factors.

As with the heat transfer results we concentrate in presenting a mean apparent friction factor,

$$f_a = -\frac{Dh}{4L} \frac{\rho}{G^2/2g_c} \left(\sum_{q_b} \left\{ p - \frac{G^2}{\rho_b g_c} \right\} \right)$$
(13)

(The local apparent friction factor, f_x , appearing in the Tables is defined in the analogous derivative form with d/dx replacing $(\frac{1}{L})$ $\stackrel{\triangle}{\circ}$). When P_b

is constant, the second term in brackets does not change and the definition reduces to that of Shah and London [7]. With constant fluid properties one solves equations (5a), (5b) and (5d) only, so the result is independent of Prandtl number and can be written as a single function $f_a(L^+)$ which approaches $f_a \cdot Re/24$ as L^+ becomes large. This function may be found tabulated in Appendix A or can be derived from earlier local results [10, 24]. For a continuous approximation, the approach of Schlünder can be used as in equation (13) to give

$$f_a \cdot Re/24 = \sqrt{1 + 0.0788/L^+}$$
 (14)

which represents the numerical results well in the immediate entry but is 4 to 5 percent high in the range, $0.05 < L^+ < 0.2$.

With varying transport properties, the energy equation (5c) is coupled to equations (5a, b and d) via the temperature-dependent viscosity and density, so the wall friction also becomes a function of the Prandtl number and the heating rate. Again the question arises as to the better method of accounting for the fluid property variations. Predictions of friction are not as well behaved as heat transfer parameters. In contrast to the heat transfer results, direct use of the average bulk properties in f: Re and L^+ does not collapse the results nicely around the prediction based on constant properties; the main effect is to spread the curves towards larger $L_{\rm ha}^+$ as Q^+ increases.

The effect of heating rate on apparent wall friction is presented in Figure 6 partially in terms of average bulk properties. That is, P_{ba} is used for the coefficient in equation (13) and Re_{ba} is defined as before but the non-dimensional length is based on inlet properties, i.e., L_{o}^{+} . With this representation heating increases f_{ba} · Re_{ba} considerably more than Nu_{ba} is raised at the same level of Q^{+} . At lengths greater than $L^{+}=0.1$ the curves with heating approach the constant properties curve only slowly, although $T_{wa}T_{ba}$ is close to unity, as the heated entry continues to affect the integrated results for downstream. Close inspection of the trends for the highest heating rates shows that as P_{c} increases a convergence – from heated entry behavior towards agreement with constant property behavior – is moved further downstream. This effect corresponds to the difference in growth of the thermal boundary layer and shear boundary layer as the Prandtl number changes: $Nu(L^{*})$ shows only a moderate effect of P_{c} , so for the same heating rate the

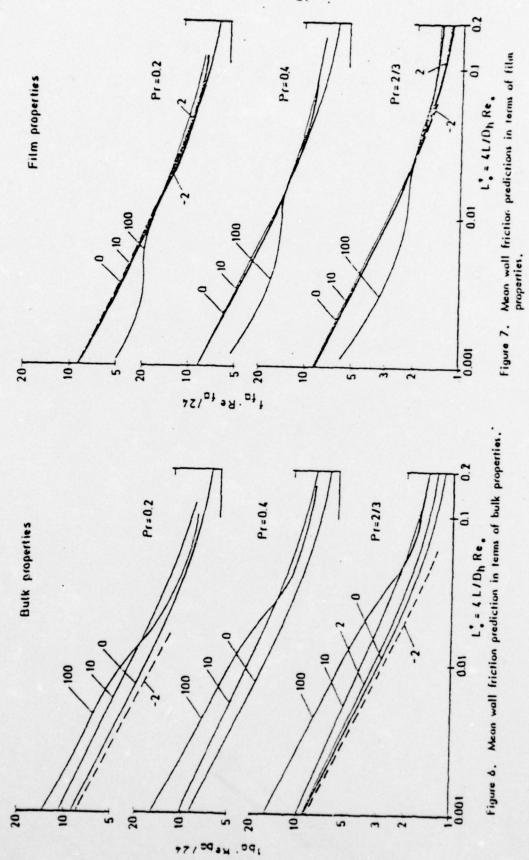
value of T_{b} is almost the same at equal value of L^{*} rather than L^{+} , thus T_{b} approaches unity for Pr = 0.2 at earlier values of L^{+} than for Pr = 2/3 and the variation of properties across the channel is less for Pr = 0.2 at the same L^{+} . While the friction predictions for heating approach the adiabatic prediction as L_{ba}^{+} increases, those for cooling diverge; this result also corresponds to the trend of property variation since the ratio T_{ba}/T_{wa} increases downstream for cooling as described earlier in the section on heat transfer.

As with the heat transfer results, the apparent effect of property variation on wall friction is sensitive to the choice of reference temperature. With average film temperature for the reference, the shape of the resulting curves differs from the shape with bulk temperature as reference. In Figure 7 the product fra · Refa/24 is plotted against Lo; P fa is used in the coefficient in equation (13) and Re_{fa} is based on μ_{fa} . There is no advantage in comparison on the basis of L_{fa} since results are shifted then further to the right (with heating) so that for L+ > 0.01 the difference from the adiabatic prediction is increased. For heating: the friction parameters are reduced for short lengths; then the predictions coverge with and cross the constant properties curve and remain slightly greater at larger distances. In comparison to the bulk property predictions, the effects for short and long ducts are approximately the same magnitude with strong heating, but for intermediate lengths and for Q+ 2 10 a display in terms of film properties shows significantly less variation. In the range $-2 < Q^+ < 2$ there is no significant effect of heating until L_0^+ approaches 0.1 with film properties and then the effect is only of the order of five to ten percent. As is the case for bulk properties, the convergence towards the adiabatic prediction is at successively greater distances (L2) as the Prandtl number increases, but for the same condition it is several times earlier with film properties. With cooling: the direction of the trends are reversed but for Q > -2 they are essentially again negligible for entry problems.

It is not clear from Figure 6 and 7 which approach is better: property ratio or film temperature. The property ratio approach would be represented as

$$(f_{ba} \cdot Re(L_o^+)) / (f_a \cdot Re(L_o^+))_{cp} = (T_{wa}/T_{ba})^q$$
 (15)

so this quotient is plotted versus temperature ratio in Figure 8 to examine the suitability of a single exponent. A complicated pattern appears. In contrast to the expectation of Kays and London [25], the general trend is a substantial



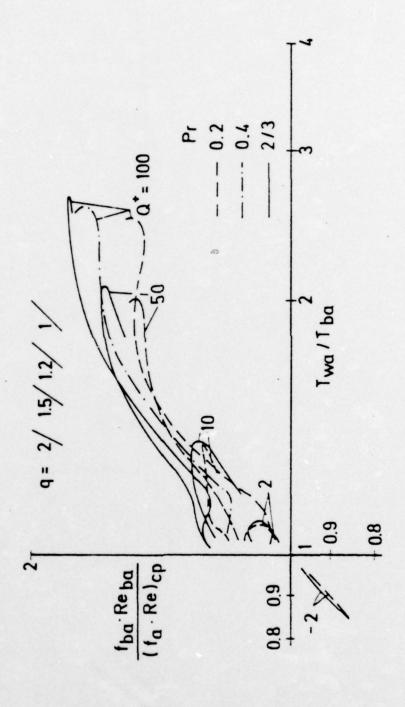


Figure 8. Examination of property ratio method for mean wall friction.

increase with temperature ratio. There are slight differences with Prandtl number but the trends are mostly the same. An exponent q of the order of unity would overpredict the friction factor at the higher temperature ratios and underpredict it at lower values. For cooling, q = 1 is valid within a few percent. For moderate heating, the necessary value of q (i.e., the slope of a line from the origin on this logarithmic plot) varies with length Lot: it is approximately constant as the temperature ratio increases with length then increases gradually as the ratio drops for successively longer ducts. The latter effect is a consequence of the slow covergence of fba · Reba to the adiabatic curve for long ducts as discussed earlier. It is seen that a function $q(L_o^+, Q^+, Pr)$ would be necessary to describe the detailed behavior. For Q = 2 and $L_0^+ < 0.6$, an exponent q = 1.5 would reduce the difference from the constant properties curve from 13 percent to a 7 percent discrepancy. With $Q^+ = 10$, q = 1.2 is a better approximation, but the discrepancy would still reach twenty percent. These comments and comparison of Figure 6 and 8 suggest that the two methods have approximately the same overall accuracy for $Q^+ \approx 2$ with a slight advantage to the film properties approach for short ducts. For moderate heating - to Q+ = 10 - the film property method is clearly superior, while at higher heating rates both methods show regions where the simple correlations would mislead the designer substantially.

It is perhaps inconvenient for the designer to have one method perform better for heat transfer while the other is preferable for wall friction, but the difficulty should be negligible provided the present definitions of the parameters are used. Once the heat transfer problem is solved for the wall temperature using average bulk properties, the average film temperature can be calculated from the results and can then be employed to predict the wall friction behavior.

Analytical correlations such as equations (12) and (14) are useful for parameter studies of systems and for initial sizing of components when hundreds to thousands of individual configurations may be calculated. When greater accuracy is needed in final design decisions – or if variable wall heat flux should be treated – the numerical analysis can be employed directly. With the direct application of the program, the question of definitions of the non-dimensional parameters is avoided; the engineer can choose definitions to suit his own convenience, including direct presentation in temperatures, pressure and lengths in units of his choice.

6. CONCLUSIONS

For heat transfer to mixtures of inert gases in short ducts formed of parallel plates, the following conclusions concerning the behavior of the mean parameters are warranted. In terms of appropriate non-dimensional variables and parameters, the effects of varying mixture concentration can be represented by variation of the inlet Prandtl number. These conclusions are based on the specific definition of parameters chosen in the present work; with strong heating rates the use of alternate definitions in the correlations can cause substantially different predictions of the heat transfer coefficient and friction factor and, consequently, of the wall temperature and pressure drop.

a) Under the constant property idealization, heat transfer and apparent wall friction parameters can be approximated by

$$Nu_{a} = \left[8.235^{2} + 1.931^{2}/(Pr^{0.254}L^{*})\right]^{1/2}$$
 and
$$(f_{a} \cdot Re/24) = \left[1 + 0.0788/L^{+}\right]^{1/2}$$

to within ten percent for L* and L+, respectively, greater than 0.001.

- b) For heat transfer in the range $-2 < Q^+ < .100$ the bulk properties/property ratio method of accounting for the effects of gas property variation with exponent p ≈ 0.005 provides better predictions than the film temperature method.
- c) For wall friction in the range $-2 < Q^+ < 10$ the film temperature method is more accurate overall than the property ratio approach.

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APPENDIX A

TABULATED RESULTS

This appendix presents reproductions of the summary tables of the computed results that were provided as output from the UNNAC 1108 computer at the Rechenzentrum, Universität Karlsruhe. All cases for which predictions were calculated are included in these tables. Space limitations prohibit listing all the predicted parameters of interest; however, the set shown should provide the basic information from which others can be calculated via their definitions and the idealized property relationships. For example,

$$L_{ba}^{\bullet} = \frac{4L}{D_{h}Re_{ba}Pr} = \frac{4L}{D_{h}Re_{o}Pr} \frac{Re_{o}}{Re_{ba}} = L_{o}^{\bullet} \left(\frac{T_{ba}}{T_{o}}\right)^{a} = \frac{4x}{D_{h}Pe_{o}} \left(\frac{T_{ba}}{T_{o}}\right)^{a}$$

The cases calculated are:

Pr	•	o ⁺
0.7 (air)		50, 0†
0.666		100, 50, 25, 10, 2, 0, -2
0.4		100, 10, 0+
0.2		100, 50*, 25*, 10, 2, 0, -0.2, -2

^{*}The case of Q^+ = 0 represents heating rates sufficiently small that the properties can be considered constant. In order to calculate these cases non-zero heat flux was chosen and the property exponents were chosen for constant properties; these are identified in the following tables by the head: VIS = .000, CON = .000, CP = .000.

^{*}Inlet pressure relatively low, leading to excessive Mach number at last few stations.

Of necessity some of the table headings are brief so the following definitions and/or comments are appropriate:

Program Symbol	Usual Nomenclature	Definition	Comment
PR, 0	Pro	(cp 4/k)0	Inlet Prandtl number
QPP.DH/KO	·T0 Q ⁺	q"Dh/koto	
GAMMA		c _p /c _v	Specific heat
VIS	a		viscosity exponent in equation (6)
CON	Ь		thermal conductivity exponent in equation (6)
СР	d		specific heat exponent in equation (6)
4X/DPEO	×.*	4x/(DhReoPr)	
TB/TI	Tbx/To · ·		local bulk temperature
TW/TB	Twx Tbx		local temperature ratio
NUBX	Nubx	hDH/k bx	local bulk Nusselt number
FR8/24	f _{bx} ·Re _{bx} /24		local apparent friction product, based on bulk properties, see pg.14
WA/BA	Twa Tba		average temperature ratio, see pg. 10
NUBA	Nuba	haDh/k ba	average bulk Nusselt number, see pg. 11
FRBA/24	fba Reba/24		mean apparent friction product, based on bulk properties, eq.(13) and pg. 15
PO-P1	Po-Px	2g_p_(pp)/C	pressure drop, note pg.9
K	= P ba (Po - PL) - 2F		$\frac{24x_{ba}^{\bullet}}{2+a} \frac{\hat{\rho}_{ba}}{(\hat{T}_{ba})^{a}} \left[\frac{\hat{T}_{bL}^{2+a}-1}{\hat{T}_{bL}-1} \right]$
			comparable to $K(x)$ of Shah and London [7] - but error in output in earlier constant property results.

LAMINAR GAS FLON BETWEEN PARALLEL PLATES

SYMMETRIC CONSTANT WALL HEAT FLUX

PR. 0= .700, QPP. OH/KO. 10= 50.0, GAMMA=1.400

PROPERTY EXPONENTS, VIS= .470.COM= .835.CP= .395

4X/DPED TB/TL T4/TB NUBS FRB/24 WA/BA NUBA FRBA/24 PO-PI 9.235 1.645 54.132 15.361 3.688-01 1 . 03-03 1.05 2.263 36.35 1.36 2.371 32.52 7.775 1.777 50.668 13.718 4.449-01 1 . 3 3 - 03 1.75-03 1.39 2.492 28.35 6.958 1.777 51.274 11.988 5.514-91 1.11 2.596 25.83 6.112 1.841 45.339 10.590 6.758-91 2.30-03 5.492 1.991 40.255 9.373 8.263-01 3.00-03 5.157 1.936 37.396 8.629 9.550-01 1.18 2.732 21.59 3.60-03 4.671 1.975 33.573 7.772 1.141+00 4.50-03 1.22 2.774 17.64 7.999 1.351+00 1.27 2.708 13.14 4.386 2.330 37.685 5.50-03 4.049 2.518 24.282 3.789 2.526 25.939 6.527 1.577+00 1.32 2.788 16.87 4.63-03 1.39 2.764 15.67 5.776 1.866+00 4.00-03 1.47 2.708 14.30 1.00-02 3.525 2.021 21.460 5.476 2.283+00 1.63 2.599 12.97 3.240 1.901 20.834 4.822 2.931+00 1 . 30-02 1 . 75-02 1.84 2.431 11.57 2.933 1.928 19.221 4.247 3.941+00 2.643 1.839 14.127 3.792 5.259+00 2.10 2.238 10.54 2 - 30 - 02 3.33-92 2.43 2.929 9.8 2.346 1.729 14.377 3.390 7.058+00 2.70 1.834 9.47 3 - 125 8 - 691+99 1.63-02 2.178 1.645 13.341 4+50-02 3.11 1.712 9.07 1.875 1.538 12.252 2.812 1.130.01 3.55 1.572 8.87 1.672 1.446 11.439 2.543 1.441+01 5.50-02 8.74 1.523 1.367 12.827 4.24 1.461 2.316 1.811+91 4 . 6)-12 1.417 1.298 13.296 4.65 1.362 8.51 2.113 2.333+91 A . 00-02 1.00-01 5.59 1.270 8 . 54 1.334 1.228 9.806 1.898 3.193.01 2.00-01 9.63 1.100 8 . 35 1.199 1.091 a . 873 1.581 1.045+02 3.00-01 13.60 1.054 8 . 3-1.123 1.050 4.612 1.457 2.312+92

1.089 1.033

1.072 1.023

4.500

A.435

8 . 22

4.00-01 17.45 1.035

5.33-91 21.22 1.024 8.27

1.397 4.267+02

1.364 7.062+02

LAMINAR GAS FLOW BETWEEN PARALLEL PLATES

SYMMETRIC CONSTANT WALL HEAT FLUX

PR.0= .700, QPP.DH/KO.TO= 200.0, GAMMA=1.400

PROPERTY EXPONENTS. VIS= .COC.CON= .CCC.CP= .CCC

4X/DPED TRITE TWITE NURY FRE/24 WAIRA NURA FREA/24

		The second secon
1-00-03	34.46 5.354	64.738 10.54R
1+15-03	32.39 5.083	40.454 9.853
1.30-03	30.68 4.760	57.294 9.287
1.50-03	28.80 4,40E	53.621 A.650
1.75-03	26.04 4.090	49.942 8.025
2.00-03	25.42 3.84A	46.972 7.512
2.30-03	23.95 3.646	44.065 7.022
2.60-03	22.73 3,441	41.673 6.620
3.00-03	21.40 3.208	39.058 6.180
3.30-03	20.56 3.041	37.415 5.902
3.60-03	19.82 2.955	35.979 5.660
4.00-03	18.98 2.830	34.322 5.384
4.50-03	18.10 2.675	32.569 5.090
5.00-03	17.36 2.555	31.085 4.542
5.50-03	16.71 2.441	29.807 4.629
6-00-03	14.15 2.350	28.693 4.442
6.40-03	15.57 2.266	27.526 4.248
7.20-03	15.06 2.173	26.508 4.079
8.00-03	19.48 2.075	25.334 3.883
9.00-03	13.87 1.978	24.095 3.676
1.00-02	13.36 1.889	23.047 3.502
1-15-02	12.73 1.792	21 - 742 3 - 283
1.30-02	12.21 1.697	20.672 3.105
1.50-02	11.66 1.605	19.507 2.910
1.75-02	11.11 1.517	18.347 2.717
2.00-02	10.AR 1.445	17.416 2.562
2.30-02	10.27 1.375	16.511 2.411
2.40-02	9.95 [-3]9	15.773 2.2A8
3.00-02	9.61 1,259	14.973 2.154
3.30-02	9.40 1.224	14.477 2.071
3.40-02	9.23 1.192	14.047 1.999
4.no-02	9.05 1.157	13.556 1.916
4.50-02	8.87 1.122	13.046 1.830
5.00-02	8.73 E-09A	12.621 1.757
5.50-02	8.63 1.076	12.263 1.696
6.00-02	8.54 1.059	11.957 1.644
6.60-02 7.20-02	8.47 1.044	-11.643 1.590
		11.376 1.544
8.00-02 9.00-02	8.36 1.023 R.32 1.014	11.077 1.492
1.00-01		10.773 1.439
1.50-01	8.29 1.010	10.526 1.397
2.00-01	8 • 25 1 • gn 2	9.773 1.266
2.50-01	8.24 1.001	9.390 1.199
3.00-01		9-160 1-160
3.50-01	8.24 1.001 8.24 1.001	9.006 1.133
4.0C-01		8.896 1.114
4.50-01	The state of the s	8.814 1.100
5.20-01		8.750 1.089
5+50-01	The second secon	8.699 1.080
3.30-01	8.24 1.001	8.657 1.073

LAMINAR GAS FLOW BETWEEN PARALIEL PLATES SYMMETRIC CONSTANT WALL HEAT FLUX

PR. 0 = .667. GPP.DH/KO.TO = 100.0. GAMMA=1.667 PROPERTY EXPONENTS. VIS= .750.CON= .750.CP= .000 4X/OPEO TB/TE TW/TB NURY FRB/24 WA/BA NURA FRRA/24 PO-PI 1.13 3.512 32.14 1.30-03 10.114 2.333 57.728 17.092 6.572-01 8 . 674 2 . 439 50 . 321 1 . 75 - 03 14.912 8.318-01 1 . 23 3 . 742 25 . 25 13.094 1.040+00 2.30-03 7.794 2.523 44.337 3.00-03 1.30 3.810 22.53 6.919 2.587 39.181 11.506 1.302+00 1.36 3.810 20.81 3 - 40 - 03 6.339 2.619 35.970 10.533 1.528+00 4.50-03 1.45 3.774 18.86 5 . ASO 2 . 641 32.367 9.432 1.864+00 5.50-03 1.55 3.696 17.26 5.410 2.638 29.398 A.552 2.247+00 1.66 3.5AA 15.93 6.60-03 5.073 2.614 26.898 7.823 2.682+00 1.80 3.443 14,42 8.00-03 4,701 2,571 24.474 7.132 3.249+00 1.00-02 2.00 3.243 13.25 4,319 2,495 21.914 6.418 4.090+00 5.677 5.442+00 1 - 30-02 2.30 2.940 11.47 3.AAA 2,366 19.211 1.75-02 2,75 2.605 10.60 3.372 2.177 16.555 4.939 7.665+00 2 - 30 - 02 3.30 2.271 9.73 2.849 1.976 14.492 4.304 1.066+01 3.00-02 9.15 4.00 1.946 2.339 1,773 12.836 3.701 1.485+01 3.40-02 4.60 1.778 8,90 2,027 1,639 11.896 3.293 1.873+01 4.50-07 5,50 1.582 A . 70 1.723 1.492 10.954 2,832 2.510+01 5.50-02 4.50 1.440 4.58 2.469 3.305+01 1.514 1.382 10.284 8.51 6.40-02 7.60 1.338 1.392 1.299 9.796 2.193 4.307.01 1.307 1.228 8.00-02 9.00 1.253 8.44 9.385 1.960 5.818+01 1.00-01 11.00 1.190 A.38 9.015 1.757 8.514+01 1.268 1.165 2.00-01 21.00 1.059 A. 28 1.197 1.056 8.334 1.529 3.728+02 3.00-01 31.00 1.030 A. 26 8 . 157 1.428 9.899+02 1,138 1,029

1-139 1-018

A . 083

1.384 2.149+03

A . 26

4.00-01 40.99 1.01A

LAMINAR GAS FLOW BETWEFN PARALIEL PLATES. SYMMETRIC CONSTANT WALL HEAT FLUX

PR. 0 - . 667. OPP. DH/KO. TO= 50.0. GAMMA=1.667

PROPERTY EXPONENTS. VIS= .750.CON= .750.CP= .000

4X/DPED TB/TI TW/TB NUBX FRB/24 WA/BA NUBA FRBA/24 PO-PI 1.06 2.428 31.39 8.494 1.736 57.544 14.514 4.478-01 1 - 30-03 1 . 75 - 03 7,477 1.810 50.165 12.781 5.591-01 1.11 2.445 24.85 6.870 1.877 44.264 11.358 6.873-01 2-30-03 6,129 1,941 39.230 10.104 8.470-01 3.00-03 1.15 2.741 22.28 5.636 1.981 36.111 1.18 2.414 20.67 3 . 40-03 9.320 9.807-01 4.50-03 1:22 2.859 18.90 5.216 2.023 32.635 8.436 1.178+00 5.50-03 1.27 2.879 17.43 4,817 2.053 29.797 7.706 1.396+00 6.40-03 1.33 2.A78 16.20 4.513 2.071 27.431 7.101 1.636+00 8.00-03 1.40 2.854 15.00 4,206 2,081 25.135 6.514 1.943+00 1 . 00-02 1.50 2.791 13,75 3,904 2,074 22.704 5.904 2.393+00 1 . 30-02 5.252 3.088.00 1.65 2.475 12.43 3.543 2.043 20.13A 1.75-02 1.87 2.494 11.15 3,154 1.974 17.591 4.609 4.184+00 2.15 2.286 10.20 2.30-02 2,803 1,877 15.554 4.079 5.622+00 3.00-02 2.50 2.059 9.51 2.440 1.756 13.851 3.605 7.608+00 2,192 1,664 12.844 3.40-02 2.80 1.902 9.16 3.289 9.427.00 1,911 1,549 11.791 4.50-02 3.25 1.71A A. A6 2.920 1.237+01 3.75 1.570 A . 49 5.50-02 1.447 1.450 11.007 2.608 1.594+01 6.40-02 4.30 1.453 8.59 1.530 1,368 10.418 2.348 2.027+01 8.00-02 5.00 1.351 A.52 1,413 1,293 9.909 2.110 2.651+01 9.44 9.440 1.00-01 6.00 1.258 1.341 1.221 1.885 3.710+01 A.30 A.540 2.00-01 11.00 1.091 1,211 1.083 1.564 1.362+02 3.00-01 16.00 1.047 A . 28 A. 287 1.436 3.275+02 1:131 1:044 4.00-01 20.99 1.029 1.098 1,028 A.27 8 - 174 1.375 6.470+02 5.00-01 25.99 1.020 1.083 1.019 A . 26 8.112 1.343 1.139+03 6.00-01 30.99 1.015 8.26 1.080 1.014 8.074 1.324 1.869+03

H

LAMINAR GAS FLOW BETWEEN PARALIEL PLATES

SYMMETRIC CONSTANT WALL HEAT FLUX

PR.C= .667, RPP.DH/KD.TC= 25.0, GAMMA=1.667

PROPERTY EXPONENTS. VIS= .750.CON= .750.CP= .000

4X/DPEC TRITI THITE NURY FR8/24 WA/BA NUBA FR84/24 1.00-03 1.03 1.686 34.88 7.883 1.347 65.189 14.181 2.823-01 1 - 30 - 03 1.03 1.758 31.15 7.218 1.385 57.678 12.618 3.356-01 1.75-03 1.04 1.840 27.58 6.450 1.429 50.279 11.164 4.135-01 2.30-03 1.06 1.921 24.61 5.832 1.473 44.381 9.909 4.988-01 3.00-03 1.08 1.997 22.08 5.188 1.517 39.350 B.833 6.034-01 3.60-03 1.09 2.051 20.46 4.848 1.548 36.251 8.138 6.868-01 4.50-03 1.11 2.108 18.72 4.463 1.584 32.806 7.391 8.108-01 6.777 9.442-01 5.50-03 1 . 14 2 . 154 17 . 29 4.094 1.614 30.008 6 - 60 - 03 1.15 2.186 16.14 3.875 1.638 27.680 6.264 1.089+00 1.20 2.210 15.02 5.770 1.272+00 8.00-03 3.678 1.660 25.437 1.25 2.224 13.84 1.00-02 3.358 1.48C 23.080 5.241 1.530+00 1 - 30 - 02 1.32 2.212 12.61 3.074 1.691 20.608 4.694 1.920+00 1.75-02 1.44 2.140 11.42 2.786 1.684 18.154 4 . 142 2 . 5 15+00 2.30-02 1.58 2.074 10.50 2.530 1.657 16.191 3.696 3.271+00 1.75 1.940 3.00-02 9.78 2.285 1.611 14.537 3.307 4.275+00 1.90 1.847 3.40-02 9.38 2.119 1.568 13.537 3.059 5.178+00 4.50-02 2.12 1.740 9.03 1.929 1.503 12.461 2.776 6.619+00 2.37 5.50-02 1 . 624 8 . 8 1 1.756 1.439 11.634 2.538 8.335+00 6.40-02 2.65 1.523 84.8 1.614 1.380 10.993 2.334 1.037+01 8.00-02 3.00 1.476 8.59 1.494 1.319 10.424 2.136 1.321+01 1.00-01 3.50 1.328 8.50 1.407 1.255 9.885 1.935 1.782+01 2.00-01 6.CC 1.130 9.33 1.237 1.112 8 . 799 1.586 5.566+01 8.50 1.071 3.00-01 8 - 29 1.144 1.064 8.463 1 - 444 1 - 218 + 02 4.00-01 11.00 1.045 8.28 1 . 374 2 . 251+02 1.102 1.042 8.306 5.00-01 13.50 1.032 8.27 8 . 217 1.080 1.030 1.335 3.747+02 6.00-01 16.00 1.024 A . 26 1.046 1.022 8.160 1.311 5.804+02

PR.0= .667, QPP.DH/x0.T0= 15.0. GAMMA=1.667

PROPERTY EXPONENTS, VIS= .750 CON= .750 CP= .000

4X/DPED TB/TI TH/TB NUBX FRR/24 HA/BA NUBA FRBA/24 PO-PI 1.00-03 1.01 1.283 34.71 6.706 1.142 65.642 12.114 2.160-01 5,987 1.159 57.558 10.758 2.529-01 5.261 1.180 50.172 9.407 3.031-01 1.01 1.316 30.94 1 . 30-03 1.02 1.356 27.27 1 . 75 - 03 1.02 1.396 24.25 4.736 1.200 44.271 8.334 3.596-01 2 - 30 - 03 7 . 415 4 . 262-01 4.265 1.222 39.245 1.03 1.437 21.72 3.00-03 3.984 1.237 34.154 6.853 4.803-01 1.04 1.467 20.15 3 . 60 - 03 3.640 1.257 37.729 6 . 221 5 . 569-01 4.50-03 1.04 1.502 18.47 5 . 707 6 . 381-01 3.387 1.274 29.954 1.05 1.534 17.05 5.50-03 3.161 1.295 27.659 5.282 7.241-01 4.60-03 1.07 1.562 15.91 1.08 1.590 14.82 4.869 8.295-01 2.940 1.376 25.452 a . no-03 4.433 9.748-01 2.717 1.324 23.143 1.00-02 1.10 1.418 13.69 2.495 1.342 20.737 3.975 1.186+00 1 . 30-02 1.13 1.444 12.53 2.269 1.357 12.371 1.75-02 1.17 1.461 11.41 3.521 1.496+00 2.085 1.364 14.490 2 - 30 - 02 1.23 1.460 10.55 3.156 1.871+00 2.844 2.351+00 1.928 1.362 14.905 3.00-02 1.30 1.641 9.84 3.60-02 1.36 1.416 9.42 1.827 1.355 12.944 2.650 2.770+00 4.50-02 1.45 1.573 1.718 1.339 12.904 2.436 3.415+00 9+11 8.84 5.50-02 1.55 1.523 1.624 1.318 12.091 2 . 261 4 . 163 + 00 4+60-02 1.66 1.472 8.72 1.546 1.295 11.449 2.116 5.027+00 1.475 1.267 15.866 1.978 6.199+00 A . 00-02 1.80 1.415 8.62 1.00-01 2.00 1.348 8.53 1.413 1.232 10.301 1.836 8.028+00 3.00 1.175 8.37 1.266 1.131 9 . 121 1.553 2.102+01 2.00-01 4.00 1.106 8 - 33 1 - 176 1 - 085 A . 720 1.429 4.114+01 3.00-01 5.00 1.n72 8.30 a . 520 1.364 7.063+01 4.00-01 1.133 1.060 A . 399 6.00 1.452 8.29 1.111 1.945 1.327 1.123.02 5.00-01 A . 320 4.00-01 7.00 1.040 8 · 2 A 1 · 101 1 · 035 1.306 1.699+02

PR.0= .667. APP.DH/KO.TO= 2.9. GAMMA=1.667 PROPERTY EXPONENTS. VIS= .750.CON= .750.CP= .000 4X/DPEO TB/TI TW/TB NUBY FRB/24 WA/BA NUBA FRBA/24 PO-PI 1.00-03 1.00 1.057 34,69 5,804 1.029 65.131 11.098 1.823-01 1.00 1.045 5.120 1.032 57.639 4.452 1.037 50.240 9.779 2 - 095 - 01 1 - 30 - 03 30.96 8.482 2.458-01 1 - 75 - 03 1.00 1.073 27.12 2.30-03 1.00 1.082 24.11 3,934 1.041 44.325 7.446 2.850-01 3.00-03 1.01 1.092 21.55 3.520 1.046 39.287 6.569 3.297-01 19.97 3 - 40 - 03 1:01 1:099 3.259 1.050 36.188 6.035 3.650-01 4.50-03 1.01 1.108 18.24 2.934 1.054 32.755 5.440 4.136-01 5.50-03 1:01 1:116 16:84 2.721 1.059 29.974 4.960 4.638-01 1.01 1.124 15.70 6.60-03 4.565 5.153-01 2.517 1.063 27.477 2.326 1.067 25.471 8.00-03 1,02 1.133 14.41 4.185 5.765-01 1.00-02 1.02 1.143 13.48 2 - 131 1 - 072 23 - 167 3.786 6.579-01 1.03 1.165 12.34 1 - 30-02 1.934 1:078 20.775 3.372 7.714-01 1.04 1.168 11.24 1.741 1.085 18.433 2.966 9.284-01 1.75-02 1.05 1.178 10.39 2.30-02 1.592 1.091 16.582 2 . 644 1 . 107+00 3.00-07 1,06 1.186 9.73 2.372 1.323+00 1.469 1.096 15.029 3 - 40-02 1 - 07 1 - 1 99 9.35 1.394 1.098 14.093 2 - 207 1 - 501 +00 4-50-02 1:09 1:191 8.99 2.027 1.762+00 1.322 1.100 13.081 5.50-02 1.11 1.190 A . 75 1.884 2.047+00 1.270 1.100 12.290 6.40-02 1-13 1-197 8.60 1.235 1.099 11.663 1.769 2.361+00 8.00-02 1,16 1.182 1,209 1,098 11.091 8.49 1 . 664 2 . 766+00 1.00-01 1 - 20 1 - 173 A.43 1 - 1 94 1 - 094 10.534 1.560 3.361.00 9.380 2 . 00-01 1.40 1.133 A . 35 1.167 1.078 1.353 6.798+00 3.00-01 1.60 1.104 A . 33 1,135 1,065 A . 983 1 . 277 1 . 099+01 1 . 235 1 . 602+01 1:113 1:055 4.00-01 1.80 1.086 9.31 8 . 77A 5.00-01 2.00 1.072 1.210 2.199+01 A.30 1.094 1.048 8 . 652

1.084 1.042

8 . 565

1.193 2.898+01

A. 29

2 . 20 1 . 0 4 1

6.00-01

PR.U= .607. UPP.DH/KU.TU= 100.0. GAMMA=1.667
PROPERTY EXPONENTS. VIS= .000.CON= .000.CP= .000

4X/DPE0	[B/T1	TW/TB	NUHX	FR8/24	WAZBA	NURA	FRBA/24	PO-PI
1-15-03			32.59	5.150		01.050	10.107	1.860-01
1.30-03			30.85	4.820		57.606		1.979-01
1.50-03			28.97	4.512		53.965		
1.75-03			27.10	4.218		50.200		2.303-01
2.00-03			25.56	3.950		47.269		2.464-01
2.30-03			24.08	3.712		44.341		
2.60-03			22.84	3.492		41.932		2.820-01
3.00-03			21.51	3.289		39.298		
3.30-03			20.06	3.157		37.642		
3.60-03			19.92	3.035		30.196	5.798	3.340-01
4.00-03			19.08	2.886		34.526	5.514	3.529-01
4.50-03			18.19	2.731		32.761	5.212	3.753-01
5.00-03			17.44	2.609		31.200	4.958	3.966-01
5.50-03			10.79	2.502		29.979	4.739	4.170-01
6.00-03			10:25	2.405		28.857	4.548	4.367-01
6.60-03			15.64	2.304		27.682	4.348	4.592-01
7.20-03			15.13	2.224		26.657		4.809-01
8.00-03			14.54	2.119		25.474		5.087-01
9.00-03			13.93	2.014		24.226		5.415-01
1.00-02			13.41	1.931		23.170		5.731-01
1.15-02			12.78	1.818		21.856		
1.30-02			12.26	1.732		20.779		
1.50-02			11.70	1.630		19.605		7.141-01
1.75-02			11.15	1.543		18.437		7.774-01
2.00-02			10.72	1.469		17.499		8.375-01
2.30-02			10.31	1.398		16.588		
2.60-02			9.98	1.340		15.844		9.718-01
3.00-02			9.64	1.279		15.039		
3.30-02			9.43	1.242		14.538		1.116+00
3.60-02			9.25	1.209		14.105		
4.00-02			9.07	1.172		13.611		1.251+00
4.50-02			8.89	1.135		13.096		
5.00-02			8.75	1.107		12.608		
5.50-02			8.64	1.085		12.307		1.520+00
6.00-02			8.55	1.068		11.997		1.606+00
7.20-02			8.48			11.681		
8.00-02			8.42	1.039		11.411		1.808+00
9.00-02				1.017				2.103+00
1.00-01			8.29	1.012		10.552		2.265+00
1.50-01			8.25	1.005		9.791		3.069+00
2.00-01			8.24	1.001		9.404		3.870+00
2.50-01			8.24	1.001		9.171		4.671+00
3.00-01			8.24	1.001		9.015		5.471+00
3.50-01			8.24	1.001		8.904	The second second	
4.00-01			8.24	1.001		8.820		7.072+00
4.50-01			8.24	1.001		8.755		7.873+00
5.00-01			5.24	1.001		8.704		8.673+00
5.50-01			8.24	1.001		8.601		9.474+00
6.00-01			8.24	1.001		8.626		1.027+01
0.50-01			8.24	1.001		8.596		1.107+01

PR. 0 = .667. QPP.DH/<0.To= -2.0. GAMMA=1.667

PROPERTY EXPONENTS, VIS= .750, CON= .750, CP= .000

AX/DED TH/TI TW/TH NUBX FRB/24 Wa/BA NUBA FRBA/24 PO-PI

5.262 1.00-03 1.00 ·942 34 · 69 .971 65 . 189 19 . 541 1 . 647 - 91 1-15-03 .938 32.57 4.898 .969 61.075 9.818 1.761-01 1.00 .967 57.690 1 - 30-03 1.00 .935 30.85 4.634 9.241 1.870-01 1 - 50 - 03 .965 53.989 . 931 28.97 4.267 A.598 2.002-01 1.00 3.955 .963 5n.284 7.956 2.155-01 1 . 75-03 1.00 · 925 27 · 02 7.440 2.298-01 . 921 25.54 3.703 .961 47.291 2.00-03 1.00 3.455 .958 44.362 6.937 2.456-01 . 916 24. 35 2.30-03 1.00 .99 .912 22.81 3.252 .956 41.951 6.524 2.603-01 2.60-03 . 99 .905 21.47 3.057 .953 30.314 6.074 2.787-01 3.00-03 . 99 2.996 5.796 2.917-01 . 902 25.62 .951 37.658 3 - 30 - 03 .99 · 498 19.82 5.549 3.039-01 2.766 .949 34.211 3 . 60 - 03 .99 .947 34.540 · 493 19 · 03 5 . 266 3 . 193-01 4.00-03 2.634 . 99 .944 37.773 4.50-03 . g88 18.14 2.494 4.766 3.375-01 .99 · #83 17 · 34 .942 31.277 4.715 3.547-01 5-00-03 2.366 · #78 16 . 73 .99 .939 29.989 4 . 498 3 . 708-01 5.50-03 2.258 .99 .874 16.17 .937 20.865 4.309 3.862-01 A+00-03 2 - 165 .99 · 869 15.57 .935 27 . 689 4.111 4.036-01 2.064 4-60-03 1.984 .99 . 864 15.04 .932 24.663 3.939 4.202-01 7-20-03 A . 00-03 . 98 · #58 14.47 1.878 .929 24.480 3.740 4.411-01 . 98 .A51 13.85 1.775 3.530 4.653-01 9-00-03 .926 24.230 1.687 . 45 13.33 .98 .923 23.174 3.353 4.882-01 1.00-02 1-15-02 1.579 . 98 · A36 12 · 62 .919 21.858 3 - 131 5 - 197-01 1 - 30 -02 .97 1.494 ·915 25 · 780 2.951 5.491-01 .AZ8 12-17 . 97 1 . 50-02 . A18 11.69 1.394 .910 19.605 2.754 5.849-01 .97 1 . 75-02 .807-11-05 .935 IA . 436 1.300 2.557 6.257-01 .96 .797 10.6n .901 17.497 2.00-02 1.226 2.400 6.628-01 .95 .787 10-19 .896 14.585 2.30-02 1.153 2.247 7.036-01 .95 .891 15.841 . 777 9.85 1.091 2 . 1 2 3 7 . 409-01 2.60-02 . 94 .765 9.50 .886 15.035 1.987 7.860-01 1.00-02 1.027 .989 .93 . 757 9.22 .883 14.533 1.903 8.173-01 3-30-02 .93 . 749 .953 9.12 .879 14.099 1.830 8.465-01 3-60-02 .92 . 740 8.99 .909 1 . 746 8 . 822-01 .875 13.603 4+00-02 .91 .729 8 . 7 . .868 .871 11.087 1.657 9.230-01 4.50-02 .90 .719 8.55 .835 .867 12.656 1.583 9.602-01 5-03-02 .863 17.292 5.50-02 .89 .737 8.41 .897 1.520 9.942-01 .782 .859 11.980 4-00-02 . 88 . 700 8.33 1.466 1.026+00 8 . 24 .87 . 687 . 757 .855 11.660 1.410 1.060+00 4+60-02 . 478 .735 .84 8.14 .852 11.387 1.362 1.091+00 7 - 20 - 02 A . 00-02 . 84 . 464 8.00 .709 .847 11.079 1.307 1.128+00 .82 .647 .681 ·841 1n.764 9.00-02 8.01 1 - 250 1 - 169+00

LAMINAR GAS FLON BETNEEN PARALIEL PLATES

SYMMETRIC CONSTANT NALL HEAT FLUX

PR.0= .400. aPP. DH/KO.TC= 100.0. GAMMA=1.667

PROPERTY EXPONENTS, VIS= .750.CON= .750.CP= .000

4X/OPED TS/TI TN/TS NURY FRR/24 NA/BA NUB4 FRBA/24 1.00-03 1.10 3.246 37.67 14.073 2.177 69.085 25.168 4.635-01 1 - 30 - 03 1 - 13 3.396 33.48 11.993 2.271 61.080 22.298 5.715-01 1.17 3.552 29.40 10.313 2.378 53.145 19.050 7.184-01 1.75-03 1.23 3.645 26.36 9.340 2.459 46.740 16.562 8.997-01 2.30-03 3.00-03 1.30 3.697 23.48 8.193 2.524 41.237 14.375 1.126+00 3.40-03 1.36 3.700 21.66 7.348 2.555 37.811 13.055 1.320+00 4.50-03 1.45 3.445 19.42 4.779 2.576 33.972 11.599 1.612+00 5.50-03 1.55 3.594 17.91 6.255 2.577 30.815 10.397 1.937+00 5.770 2.560 28.171 9.442 2.304+00 6.40-03 1.66 3.501 14.49 5.445 2.522 25.593 8 . 00 - 03 1.80 3.347 15.11 8.533 2.781+00 1 . 00 - 02 2.00 3.173 13.67 4.989 2.452 22.873 7.625 3.485+00 1 - 30 - 02 2.30 2.907 12.23 4.549 2.329 20.006 6.714 4.610+00 1.75-02 5.839 6.447+00 2.75 2.544 10.90 4.020 2.146 17.194 2.30-02 5.109 8.913+00 3.30 2.234 10.00 3.425 1.950 15.015 3.00-02 4.00 1.941 9.19 2.811 1.753 13.272 4.411 1.232+01 3.40-02 9.12 4.60 1.759 2.416 1.623 12.284 3.931 1.545+01 5.50 1.570 4.50-02 A. 49 2.020 1.482 11.290 3.373 2.049+01 5.50-02 4.50 1.432 8.74 1.734 1.375 10.575 2.919 2.462+01 6 - 60 - 02 7.60 1.333 4,43 1.547 1.294 10.049 2.559 3.410+01 8.00-02 9.00 1 . 751 4.53 1 . 415 1 . 226 9.598 2 . 246 4 . 498+01 1.00-01 11.00 1.178 4.45 1 - 340 1 - 143 9.185 1 . 964 6 . 369 + 01 2.00-01 21.00 1.058 9.31 1 - 214 1 - 956 8 . 414 1.598 2.474+02 3.00-01 31.00 1.030 A . 28 1 - 137 1 - 029 8.205 1.460 6.194+02 40.99 1.399 1.273+03 4.00-01 9.27 1.719 1.113 1.018 8 . 114 5.00-01 50.99 1.012 8.063 1.370 2.363+03 9.26 1.119 1.012 2.28 6.09-01 60.98 1.009 1 - 297 1 - 009 1.367 4.291+03 A . n 3 3

PR. 0= .400, QPP. DH/20. TO= 10.0. GAMMA=1.667

PROPERTY EXPONENTS, VIS# .750, CON# .750, CP# .000

4X/DRED TB/TI TW/TB NUBX FRR/24 WA/RA NUBA FRRA/24 1.00-03 1.01 1.267 35.84 8.607 1.134 60.301 15.603 1.715-01 7.688 1.150 61.290 13.875 2.316-01 1 . 30 - 03 1.01 1.298 32.77 1.75-93 6.804 1.170 53.377 12.135 2.425-01 1.02 1.337 28.74 10.744 2.887-01 2.30-03 1.02 1.376 25.59 6.048 1.190 47.051 5.482 1.211 47.660 9.545 3.429-01 1.03 1.416 22.85 3.00-03 1.04 1.444 21.14 5. 39 1.226 39.343 A.813 3.871-01 3.63-03 1.04 1.479 14.31 4.620 1.245 34.667 7.989 4.498-01 4+50-03 1.06 1.510 17.84 4.247 1.262 31.639 7.317 5.164-01 5.50-03 3.942 1.278 20.226 6.753 5.863-01 4.60-03 1.07 1.538 16.61 3.663 1.294 24.857 6 . 207 6 . 721-01 4.00-03 1.58 1.566 15.45 5.629 7.905-01 1.10 1.494 14.24 3.356 1.311 24.379 1.00-02 1 - 30 - 02 5.021 9.621-01 3.048 1.330 21.800 1.13 1.421 13.60 1.75-02 1.17 1.639 11.85 2.743 1.345 19.262 4.415 1.212+00 2.497 1.353 17.247 3.927 1.514+00 1.23 1.440 10.88 2 . 30-02 3.509 1.897+00 1.30 1.423 10.14 2.285 1.352 15.546 3.00-02 1.60-02 9.72 2.150 1.346 14.517 1.36 1.400 3 . 250 2 . 228 +00 4 . 5:, -02 1.45 1.500 9.31 2.003 1.331 11.400 2.962 2.735+00 1.55 1.513 9.34 1.875 1.312 12.525 2.727 3.317+60 4.50-02 A+60-02 1.66 1.463 8.89 1.770 1.229 11.833 2.532 3.981+00 8.74 2.344 4.869+00 1.80 1.408 1.667 1.242 11.203 A . 60-02 2.00 1.343 8.64 1.569 1.229 1: .589 2.147 6.229+00 100-01 2.00-01 9.290 3.50 1.173 8.44 1.335 1.130 1.721 1.529+01 1.00-01 4.00 1.106 8.34 1.215 1.595 A.837 1.535 2.853+01 4.00-01 5.00 1.072 8.33 1.157 1.060 1.438 4.717+01 A . 607 K+00-01 6.00 1.052 8.37 1.125 1.045 4.468 1 . 382 7 . 257+91 A.00-01 7.00 1.40 8.3 1.106 1.035 0.375 1.347 1.063+02

LAMINAR GAS FLOW BETWEEN PARALIEL PLATES

SYMMETRIC CONSTANT NALL HEAT FLUX

PR.D. .400. GPP.DH/KD.TO= 100.0. GAMMA=1.667

PROPERTY EXPONENTS. VIS= .000.CON= .000.CP= .000

PROPERTY EXPONE	N15. Y(3= .UUU.CU	N= 4000.CP= 4000
WANDPED TRATE TAKE	TB NURY FRR/24 A	AZBA NUBA FRBAZZ4 PO-PI
1.00-03	36.29 7.122	69.499 13.878 1.332-01
1.15-03	14.64 6.535	65.099 12.957 1.430-01
1.30-03	32.77 6.164	61.476 12.184 1.521-01
1.50-03	30.75 5.840	57.513 11.365 1.637-01
1.75-03	24.73 5.343	53-545 10-539 1-771-01
2.00-03	27.08 5.042	50.340 9.870 1.895-01
2.30-03	25.48 4.750	47.202 9.224 2,037-01
2.40-03	24-16 4-448	44.619 A.69R 2.169-01
3.00-03	22.72 4.190	41.795 8.099 2.333-01
3.30-03	21.91 4.045	40.020 7.737 2.451-01
3.40-03	21.02 3.849	38.469 7.424 2.566-01
4.00-03	20.10 3.617	36.678 7.049 2.707-01
4.50-03	19.15 3.448	34.783 6.659 2.877-01
5.00-03	18.34 3.289	33.179 6.329 3.038-01
5.50-03	17.64 3.160	31.799 6.046 3.193-01
6.00-03	17.03 3.029	30.593 5.301 3.341-01
6.60-03	16.40 2.899	29.331 5.541 3.511-01
7.20-03	15.45 2.780	28.231 5.317 3.675-01
8.00-03	15.22 2.637	26.961 5.054 3.881-01
9.00-03	14.56 2.511	25.62n 4.779 4.129-n1 24.487 4.545 4.363-n1
1.00-02	14.01 2.392	
1.15-02	13.32 2.252	23.075 4.254 4.697-01
1.30-62	12.77 2.140	21.918 4.016 5.012-01
1.50-02	12:17 2:012	20.458 3.757 5.410-01
2.00-02	11.58 1.888	19.402 3.497 5.875-01 18.395 3.290 4.314-01
2+30-02	11.11 1.792	17.415 3.087 6.816-01
2.40-02	10.30 1.614	16.614 2.921 7.291-01
3.00-02	7.93 1.533	15.747 2.741 7.893-01
3.30-02	9.70 1.482	15.208 2.629 8.327-01
3.60-02	9.50 1.433	14.740 2.531 8.747-01
4.00-02	9.29 1.378	14.206 2.419 7.285-01
4.50-02	9.09 1.323	13.449 2.299 9.932-01
5.00-02	8.93 1.27A	13.185 2.199 1.056+00
5.50-02	A.AD 1.240	12.792 2.113 1.116+00
6.00-02	8.70 1.20A	12.455 2.039 1.175+00
6.40-02	A.AD 1-175	12-109 1-962 1-243+00
7.20-02	A.53 1.147	11.414 1.895 1.310+00
8.00-02	9.45 1.117	11.482 1.819 1.397+09
9.00-02	8.39 1.089	11.142 1.739 1.503+00
1.00-01	9.35 1.069	10.865 1.673 1.606+00
1.50-01	8.27 1.022	10.013 1.440 2.102+00
2.00-01	8.25 1.007	9.574 1.347 2.587.00
2.50-01	8.24 1.003	9.302 1.279 3.069+00
3.00-01	2.24 1.001	9.130 1.233 3.550+00
3.50-01	9.24 1.001	9.002 1.199 4.030+00
4.00-01	9.24 1.001	8.907 1.175 4.510+00
4-50-01	8.24 1.001	8.432 1.155 4.991+00
5.00-01	8.24 1.001	8.773 1.140 5.471.00
5.50-01	8.24 1.001	8.724 1.127 5.951+00
6.00-01	8.24 1.001	8.484 1.117 4.432+00
6.50-01	8.24 1.001	8.449 1.108 6.912+00

LAMINAR GAS FLOW BETWEEN PARALLEL PLATES

SYMMETRIC CONSTANT WALL HEAT FLUX

PR.O= .200. GPP.DH/KO.TO= 100.0. GAMMA=1.667

PROPERTY EXPONENTS. VIS= .750.CON= .750.CP= .000

4X/DPED	TB/TI	TWITE	NUBX	FR8/24	MAZBA	NURA	FRBA/24	K
1.00-03		3.120				74 - 157		. 258
1.30-03	1 - 13	3.245	35.80	16.944	2.200	65.471	30.206	• 310
1.75-03	****	3.410	31.45	13.839	2.300	56.A7A	25.992	.381
2+30-03	1 . 23	3.504	27.91	11.789	2.380	49.944	22.394	.460
3.00-03	1.30	3.542	24.76	9.697	2.447	43.987	19.104	.548
3.60-03	1.36	3.548	22.80	R.716	2.479	40.273	17.132	.619
4.50-03	1.45	3.540	20.59	7.925	2.503	36 - 115	14.914	.716
5.50-03	1.55	3 - 475	18.78	7.377	2.504	32.700	13.222	.817
6.60-03	1.66	3.394	17.24	6.488	2 . 493	29 . A54		.920
8.00-03	1.80	3.245	15.79	6.503	2.456	27.069	10.605	1 . 045
1.00-02	2.00	3 . OAS	14.26	6.045	2.390	24 - 139	manufacture (see) and (see) and	1.212
1.30-02	2.30	2.832	12.72	5.716	2.276	21.071	8.279	1 . 445
1.75-02	2.75	2.504	11.33	5 . 1 . 1	2.103	18.061	7.274	1.771
2.30-02	3.30	2.192	10.39	4.506	1.915	15.736		2 . 1 28
3.00-02	4.00	1.905	9.77	3.705	1 . 724	13.886	5.655	2.510
3.40-02	4.60	1.728	9.50	3.141	1.598	12.841		2.774
4.50-02	5.50	1.547	9.27	2.610	1.462	11.794	And the second s	3.087
5.50-02	6-50	1.416	9.08	2.203	1.361	11.038	the problem of the second seco	3.351
6-60-02	7.60	1 . 322	8.92	1.915	1 . 285	10.473	CONTRACTOR OF STREET STREET, S	3.570
8.00-02	9.00	1.744	9.76	1 . 687	1.220	9.975		3 . 781
1.00-01	11.00	1 - 175	A . 41	1.524	1.160	9.505	2.418	4.003
2 . 00-01	21.00	1.058	8.36	1.266	1.055	8 . 581	1.764	5 . 248
3-00-01	31.00	1.030	A. 71	1.158	1.029	8.308	1.545	6.175
4.00-01	40.99	1.018	A . 28	1.116	1.018	8 . 186	and the state of t	7.367
5.00-01	50.99	1.012	A . 25	1.097	1.012	A.116		9.286
6.00-01	60.98	1.009	8.28	1.094	1.009	8 - 072		12.607

PR:G= +200: QPP-DH/KO.TO= 50.0: GAMMA=1-667

PROPERTY EXPONENTS, VIS= +750, CON= +750, CP= +000

4X/DPEO	18/11	TW/TB	NUBX	FRB/24	WA/8A	NUBA	FR8A/24	K
1.00-03	1.05	2 . 152	39.84	17-196	1.590	74.627	29.429	• 188
2.00-03	1.10	2.437	29.45	12.322	1.753	53.831	21.622	•301
3.60-03	1.18	2.636	22.87	8.984	1.884	40.889	16.240	. 446
6-00-03	1.30	2.715	18.42	7.002	1.969	32.182	12.438	.627
1.00-02	1.50	2.652	14.89			25 . 283		.881
2.00-02	2.00	2.301	11.42	4.508	1.867	18.194	6.705	1.407
3.60-02	2.80	1.844	9.77			13.960		2.076
6.00-02	4.00	1 . 482	9 - 17			11.546	3.801	2.769
1.00-01	6.00	1 . 248	8.77	1.741	1.212	10.050	2.734	3.452
3.50-01	18.50	1.036	8.35			8.411	1.605	7-107

LAMINAR GAS FLOW BETWEEN PARALLEL PLATES

SYMMETRIC CONSTANT WALL HEAT FLUX

PR.0= .200. QPP.DH/KO.TO= 25.0. GAMMA=1.667

PROPERTY EXPONENTS. VIS= .750.CON= .750.CP= .000

			ANTO CONTRACTOR OF					
4X/DPE0	TB/T1	TWITE	NUBX	FRB/24	WA/BA	NUBA	FRBA/24	- K
1-00-03	1.03	1+602	39.80	14.757	1.305	75.001	25.273	•142
2-00-03				10.987				• 223
3-60-03	1.09	1.945	22.75	8.359	1.493	41.212	14.456	.326
6-00-03	1.15	2.062	18.44	6.570	1.568	32.591	11.407	+456
1.00-02		2 . 121				25.843	8 . 9 1 4	.637
2.00-02		2+044		And the second of the second o		18.977	6.359	1.008
3-60-02		1.813				14.787	4.814	1.488
4.00-02		1 . 5 4 3				12.234	3.759	2.052
1.00-01		1.314				10.565	2.867	2.709
3.50-01				1.218			1.640	4.868

Inlet pressure was relatively low on these two calculations, leading to Mach number greater than 0.2 for $x^* > 0.3$.

PR.0= +200, GPP.DH/+0.TO= 10.0, GAMMA=1.667

PROPERTY EXPONENTS. VIS= .750, CON= .750, CP= .000

PO-PI 4X/DPED TB/TI TW/TB NUBX FRB/24 W4/RA NUBA FRBA/24 1.00-03 1.01 1.246 39.89 12.320 1.124 75.346 21.958 1.266-01 1.01 1.276 35.41 10.974 1.139 64.591 19.518 1.495-01 1-30-03 9.607 1.158 57.935 17.107 1.813-01 1.75-03 1.02 1.313 31.02 1.02 1.350 27.49 8.633 1.177 51.009 15.152 2.171-01 2.30-03 1.03 1.388 24.50 7.628 1.197 45.105 13.465 2.596-01 3.00-03 1.04 1.415 22.64 7.089 1.211 41.470 12.418 2.940-01 3.60-03 1.35 1.449 20.69 6.368 1.229 37.441 11.243 3.433-01 4.50-03 1.06 1.479 18.90 5.895 1.246 34.176 10.276 3.954-01 5.50-03 1.07 1.507 17.65 5.403 1.261 31.477 9.464 4.503-01 4-60-03 1.08 1.534 16.37 4.980 1.277 2a.880 A.676 5.181-01 A . 00-03 1.10 1.562 15.04 4.519 1.294 24.165 7.835 6.111-01 1-00-02 4.056 1.313 23.340 6.944 7.456-01 1.13 1.589 13.70 1 - 30-02 1.17 1.408 12.40 3.602 1.329 20.561 6.056 9.415-01 1.75-02 3.244 1.337 14.355 5.340 1.176+00 1.23 1.A11 11.39 2.30-32 3.00-02 1.30 1.597 13.54 2.937 1.337 14.496 4 . 727 1 . 474 - 00 1.36 1.576 19.13 2.753 1.332 15.369 4.348 1.729+00 3.60-02 4.50-02 1.45 1.538 9.60 2.546 1.319 14.148 3.930 2.118.00 9.40 1.55 1.494 2.371 1.390 13.191 5.50-02 3.591 2.560+00 1.66 1.447 9.21 2.223 1.279 12.434 4-60-02 3.310 3.060+00 2.377 1.254 11.742 1.86 1.375 A.0J-02 9.04 3.040 3.721+00 2.00 1.333 8.91 1.925 1.222 11.066 2.757 4.716+00 1.00-01 3.00 1.170 8.60 2.088 1.085+01 2.00-01 1.514 1.128 9.608 4.00 1-104 8.44 1.319 1.084 0.071 1.782 1.903+01 3.00-01 1.225 1.059 9.787 4.00-01 5.00 1.071 8.40 1.615 2.983+01 5.00-01 6.00 1.752 8.34 1.172 1.045 8.612 1.515 4.384.01 4.00-01 7.00 1.040 8.31 1.140 1.035 0.493 1.450 6.168+01

PR.0= +200, QPP.DH/y0.Tn= 2.0, GAHMA=1+667

PROPERTY EXPONENTS, VIG. .750.CON= .750.CP= .600

4X/DPFC TB/TI TW/TB NUBX FRR/24 HA/RA NUBA FRBA/24 PC-PI 1.00-03 1.00 1.050 49.02 10.401 1.025 76.594 19.912 9.996-02 9.147 1.028 64.832 17.585 1.154-01 1.00 1.056 35.51 1 . 30-03 1.00 1.464 31.04 9.065 1.032 58.165 15.259 1.358-01 1 . 75 - 03 7.860 1.936 51.228 13.397 1.580-01 1.00 1.072 27.50 2.30-03 1.01 1.081 24.46 6.285 1.041 45.310 11.820 1.835-01 1.00-03 5.777 1.044 41.669 10.837 2.632-01 1.01 1.087 22.6-3.60-03 4.5u-g3 1.01 1.096 20.57 5.214 1.048 37.633 9.768 2.311-01 1.01 1.164 18.92 4,734 1.052 34.365 8.888 2.594-61 5.50-03 A.60-03 1.01 1.111 17.57 4.389 1.056 31.663 8.156 2.883-01 7.456 3.230-01 A.00-03 1.02 1.119 16.2a 4.003 1.060 29.067 1.00-02 1.02 1.129 14.94 3.645 1.065 24.356 6.715 3.688-01 3.239 1.071 23.541 5.941 4.321-01 1.30-02 1.03 1.141 13.6% 1.75-02 1.03 1.153 12.34 2.861 1.078 20.782 5.174 5.191-01 1.65 1.164 11.32 2.559 1.684 18.600 4.564 6.177-01 2 - 34 - 02 1.00-02 1.06 1.172 10.52 2.308 1.089 14.767 4 • 6 45 7 • 351 = 01 1.07 1.176 10.04 2,153 1.091 15.658 3 . 726 8 . 310 - 01 3 . 60 - 02 1.69 1.179 9.58 1.981 1.694 14.455 3.373 9.695-01 4.50-02 1.849 1.095 13.509 3.088 1.118+00 1.11 1.180 9.27 5.50-02 1.13 1.178 9.65 1.743 1.094 12.754 2.853 1.278+06 A+60-02 1.16 1.174 8.80 1,647 1.093 12.058 2.629 1.479+00 8.00-02 8.74 1,559 1.091 11.370 2.397 1.765+00 1.00-01 1.20 1.166 1.370 1.076 9.902 1.40 1.130 8.54 1.862 3.272*00 5.00-01 9.369 1 . 644 4 . 972+00 3.00-01 1.60 1.104 8.44 1.283 1.064 1.80 1.65 9.087 1.523 6.913+00 8.42 1.230 1.055 4.00-01 A. 910 1 . 446 9 . 1 26+00 2.00 1.071 8.30 1.193 1.547 5.06-01 A . 787 1.393 1.163+01 2.20 1.000 2.37 1.165 1.241 A.CC-01

LAMINAR GAS FLOW BETWEEN PARALLEL PLATES

SYMMETRIC CONSTANT WALL HEAT FLUX

PR. 1= .200, QPP. DH/KD. TO= 100.0, GAMMA=1.667

PROPERTY EXPONENTS, VIS= ,003, CON= .000, CP= .000

4X/0PED	+8/TI	TWATE	мия х	F 38/24	VA /BA	WIIBA F	RBA/2
17/ - 12/		1.17.13	100				
1.30-03			49.97	17.283		75.657	19.417
1.15-03			37.50			75.853	18.118
1.30-03			35.54	9.847		64.897	17.092
1.75-03			33.37	7.566		5a.23g	
The second second second			31.00	7.198		54.725	14.771
2.03-03			29 · 2n	6.639		51.291	13.841
2.33-03			27.52	The state of the s			12.179
2.53-03			26 - 97			42.465	
3.30-03			24.40	5.688		45.373	11.343
1.30-03			23.47	5.353		41.731	10.401
3.61-03			22.62	5.093		39.779	9.882
4.50-03			21.69	4.768		37.695	9.333
5.00-03			19.50	4.535		35.937	8.859
5.50-03			18.9	4.394		34.424	8.461
4.00-03			18.25	4.211		37.104	8.113
4.60-03			17.54	4.032		31.722	7.750
7-20-03			16.95	and a second of the last of the second of th		30.516	
A • 7) - 73			15.24	3.615		20.125	7.057
9.03-03			15.54	3.451		27.656	6.567
1.00-02			14.93	3.293		24.413	6.336
1.15-02			14.1A	3.575		24.866	5.92
1.31-02			13.54	2.995		23.598	5.58
1.50-02			12.90	2.714		22.216	5.211
1.75-02			12.25	2.535		27.839	4.840
2.00-02			11.71	2.395		19.733	4.542
2.30-02			11.24	2.260		14.657	4.25
2.60-02			10.84	2.142		17.778	4.015
3.00-02			10.42	2.019		16.825	3.75
3.30-92			10.17	1.945		14.232	3.599
3.60-02			9.95	1.879		15.718	3.454
4.00-02			9.77	1.795		15.129	3.29
4.50-02			9.40	1.7.79		14.515	3.12
5.00-02			9.29	1.641		14.902	2.976
5.53-02			9.14	1.584		11.567	2.852
4.99-92			9.72	1.533		13.193	2.744
6.60-02			A.97	1.478		12.809	2.631
7-23-02			A . 8 .	1.432		12.479	2.533
8.00-02			8.71	1.377		17.107	2.420
9.00-02			8.62	1.322		11.725	2.301
1.00-01			8.55	1.278		11.411	2.200
1.50-01			8.40	1.140		17.432	1.861
2-00-01			8.37	1.074		9.915	1.670
2.53-01			A . Za	1-039		2.592	1.546
3.00-01			8.24	1.021		7.371	1.46
3.50-01			8 . 25	1.011		0.212	1.396
4.33-01			8 . 24	1.906		0.091	1.346
4.50-01			A . 24	1.004		2.997	1.310
5.00-01			8.24	1.072		n . 921	1.279
5+50-01			A . 24	1.202		4.859	1.254
6.00-01			9.24	1.001		a . 307	1.233
4.53-31			8 . 24	1.001		3.764	1.215

PR.U= .200. UPP.DH/KU.TU= -.2. GAMMA=1.667

PROPERTY EXPONENTS. VIS= .750.CON= .750.CP= .000

4X/UPEU TB/TI TN/TB MURX FPB/24 WA/BA MURA FPRA/24 PO-PI .498 75.494 19.349 9.265-00 9.807 1.00-03 .995 40.04 1.00 8.700 .997 06.787 17.029 1.059-01 1.30-03 . 494 35.03 1.00 7.566 .997 50.155 14.798 1.231-01 .994 31.07 1.75-05 1.00 0.627 2.30-05 1.00 .93 27.60 .996 51.240 12.875 1.414-01 5.787 .990 45.351 11.306 1.01d-01 3.00-03 1.00 . 492 24.52 5.309 .996 41.720 10.350 1.776-01 3.00-03 1.00 .491 22.00 4.777 . 495 37.689 4.50-03 1.00 . 490 20.61 9.281 1.989-01 8.421 2.204-01 4.355 5.50-03 1.00 .409 10.91 . 495 34 . 425 .994 31.729 7.706 2.417-01 6.00-03 1.00 .989 17.50 3.965 3.592 . 494 29 . 132 7.021 2.606-01 8.00-03 1.00 . 408 16.20 .987 14.94 3.225 .995 20.426 6.203 2.982-01 1.00-02 1.00 1.30-02 1.00 .905 13.57 2.871 .995 23.611 5.543 3.400-01 . 984 12.24 2.502 . 992 20 . 850 4.804 3.902-01 1.75-02 1.00 .491 18.607 4.219 4.557-01 2.30-02 1.00 .402 11.23 2.223 . 44 3.724 5.227-01 .981 10.40 1.983 .990 10.934 3.00-02 .99 4.95 1.044 .990 15.725 3.424 5.748-01 3.00-02 . 401) .979 1.678 .969 14.521 . 99 3.001 0.459-01 4.50-02 9.40 .99 .989 13.573 2.824 7.180-01 . 978 1.550 5.50-02 9.12 . 44 1.445 .977 .989 12.314 2.604 7.910-01 6.00-02 0.50 8.00-05 .90 .410 0.00 1.343 .968 12.111 2.394 8.767-01 .90 .970 1.242 .986 11.413 2.176 9.887-01 1.00-01 8.52 .974 1.023 .987 9.910 2.00-01 . 40 8.27 1.045 1.448+00 3.00-01 .94 . 973 .962 .987 9.360 1.435 1.841+00 0.21

LAMINAR GAS FLOW BETWEEN PARALLEL PLATES SYMMETRIC CONSTANT WALL HEAT FLUX

PR,0= .200, QPP.DH/KO.TO= .2.0, GAMMA=1.667

PROPERTY EXPONENTS, VIS= .750, CON= .750, CP= .000 4X/DPED TB/TI TW/TB NUBX FRA/24 WA/BA NUBA FRBA/24 .975 75.714 18.807 8.629-02 1.00-03 1.00 . 950 40.11 9.455 .973 70.912 17.554 9.229-02 1 - 15 - 03 .947 37.63 8.824 1.00 . 944 35.57 .972 64.956 16.507 9.779-02 1 . 30-03 1.00 8.226 .970 67.627 15.364 1.045-01 1 - 50 - 03 1.00 .940 33.35 7.642 .935 31.12 7.081 .968 5a.290 14.215 1.123-01 1 . 75 - 03 1.00 7 . 00 - 03 1.00 . 931 29.31 6.668 .966 54.785 13.299 1.195-01 6.213 2-30-03 .963 51.352 12.407 1.275-01 1.00 . 027 27.54 . 923 26 . 09 .961 44.526 7 . 60 - 03 .99 5.796 11.669 1.348-01 .959 45.435 .99 . 918 24.51 5.412 10.855 1.438-01 3.00-03 .99 .957 43.491 .914 23.51 10.355 1.502-01 5 - 164 3 - 30 - 03 .99 4.870 .955 41.793 9.912 1.562-01 .910 22.67 3 - 60 - 03 .99 4.584 .953 39.832 9.393 1.634-01 4.00-03 .906 21.69 .99 .951 37.756 4.314 8.848 1.720-01 4.50-03 .901 20.57 .99 · 897 19 · 6 8 4.142 .949 35.999 A . 384 1 . 799-01 5.00-03 3.953 .946 34.487 5.50-03 .99 · 892 18.97 7 . 999 1 . 875-01 · #88 18 · 24 3.698 .99 .944 32.167 7 . 652 1 . 945-01 4.00-03 .99 3,541 .942 31.785 7 . 286 2 . 021 - 01 . A83 17.55 4.60-03 . 99 6.974 2.096-01 · A79 16.94 3.385 .940 30.579 7 - 20 - 03 . 98 . A73 16 . 25 .937 29.189 6.605 2.185-01 3,177 H+00-03 .98 . 467 .934 27.719 15.57 3.002 6 - 221 2 - 290-01 9.00-03 . 98 . 461 14.90 .931 24.476 2,832 5.895 2.385-01 1.00-02 2.637 .927 24.929 1 - 15-02 . 98 · a53 14 · 15 5.490 2.514-01 .97 . 445 13.52 2.467 .924 23.660 5.156 2.630-01 1 - 30-02 .97 2.292 4.793 2.767-01 . A 3 6 12.85 .919 27.27A 1 . 50-02 .96 .914 2n. 900 4.432 2.916-01 ·A25 12 · 20 1 . 75-02 2-120 1.972 .95 19.794 .910 4.143 3.045-01 2.00-02 . 416 11.67 3.859 3.177-01 . 95 1,828 19.71A . 805 11.14 . 905 7.30-02 1.716 3 . 628 3 . 291-01 . 95 . 796 10.75 17.838 2.60-02 .901 1,597 . 94 . 784 10.32 .895 14.884 3.376 3.416-01 3.00-02 .93 14.289 3 . 220 3 . 497-01 3 - 30 - 02 .776 10.05 1.529 .892 1,454 .888 15.773 . 93 . 768 3.085 3.568-01 3 - 60 - 02 9.87 1.360 .92 . 758 9.54 .884 15.182 2.928 3.641-01 4.00-02 . 91 9.31 1 - 276 . 747 .879 14.563 2.762 3.711-01 4.50-02 9.10 2.622 3.760-01 .90 1.206 .875 14.046 . 736 5.00-02 8.97 .89 .725 1.139 2.503 3.792-01 .871 13.606 5.50-02 . 88 8 . 77 1.078 .867 13.227 2.398 3.807-01 4.00-02 .715 .862 12.835 .87 8.67 1.017 2.290 3.804-01 4 - 60-02 .733 .858 17.497 .956 2 . 195 3 . 784-01 .86 . 491 8.50 7 . 20 - 02 .879 . 675 8.34 . 84 2.085 3.727-01 A+00-02 .852 17.114 8 . 21 1 . 968 3 . 615-01 .82 . 455 .792 .845 11.715 9.00-02

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